

Computational Studies on the Mechanism of Alkane Dehydrogenation on 2D layered TiS₂ and the Electronic Structure of a Manganese Dimer

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Abstract

Firstly, the possibility of dehydrogenating linear and cyclic alkanes catalyzed by 2D layered Transition Metal Dichalcogenide, 1T-TiS₂, was investigated with Density Functional Theory. The dehydrogenation mechanisms for cyclohexane and *n*-butane as representative reactants, were explored. The most likely reaction mechanism for C–H was found to occur on an edge-based S–S couple for both C–H activation steps. The first C–H activation reaction exhibited a reaction energy of –14.2, –12.7 and –12.0 kcal/mol for cyclohexane and *n*-butane addition, respectively. In contrast, the second C–H activation on the neighboring C-atom, was found to be comparatively unfavorable, resulting in an increase in energy to 19.6, 22.5 and 19.0 kcal/mol for formation of cyclohexene, 1-butene and 2-butene, respectively. Since there have been no published studies on the dehydrogenation of linear and cyclic alkanes on TiS₂ thus far, this study may provide a guideline to further computational or experimental investigations.

Secondly, the electronic structure of a dimeric manganese hydride catalyst supported by β -diketiminato ligands, [(2,6-*i*Pr₂PhBDI)Mn(μ -H)]₂, was investigated with Density Functional Theory. An apparent triple bond between two manganese centers was anticipated to compensate for the electron-deficient nature of each metal center. However, our calculations interestingly revealed the absence of a multiple bond between the metal centers. In accordance with experimentally determined Heisenberg exchange coupling constants of –10.2 cm⁻¹, the calculated J_0 value of –10.9 cm⁻¹ confirmed that the ground state involves antiferromagnetic coupling between high spin *ds* Mn(II) centers. The effect of steric bulk on the bond order was interrogated via a model study with the least sterically demanding version of the β -diketiminato ligand and was found to be negligible. Mixing between metal- and ligand-based orbitals was alternatively designated as the main cause of the absence of a metal–metal multiple bond. Moreover, the specificity of hydrides providing only *s*-orbitals affords a relatively close positioning of the metal centers, while bridging ligands including *p*-orbitals perturb a bonding orbital between metal centers, lengthening the Mn–Mn distance. The proximity of the metal centers in [(2,6-*i*Pr₂PhBDI)Mn(μ -H)]₂ leads to an increase in Pauli repulsion, resulting in destabilization of the dimer. The accessibility of the monomeric species may be the origin of the catalytic activity that [(2,6-*i*Pr₂PhBDI)Mn(μ -H)]₂ exhibits.

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Introduction

This thesis describes the efforts and results of two different projects that are based on computational research. The first is a completely theoretical study on the dehydrogenation of linear and cyclic alkanes catalyzed by 2D layered TiS_2 , which was investigated by the author of this thesis alone. The second is a collaborative effort between the author of this paper, Changjin Oh, a PhD student at Professor Baik's group, with who the computational part of this research was carried out, Professor Ryan Trovitch from Arizona State University, and Professor John Anderson from The University of Chicago, both of whom performed experiments and measurements supporting the computational investigation. This project concerns the electronic structure of a β -diketiminato manganese dimer with bridging hydrides, which Professor Trovitch and coworkers previously synthesized and used for the first time to catalyze hydrosilylation of alkenes.¹ This research is currently being prepared for publication.

The purpose of this division is to show the different applications of Density Functional Theory and computational modeling. The first investigation is a purely theoretical study to gain insight into the possibilities of performing a certain catalytic reaction and elucidating the potential reaction mechanism. On the other hand, the second is an investigation into the electronic structure of an organometallic compound in order to better understand the origins of the catalytic activity that the compound in question exhibits.

Density Functional Theory

In order to use molecular quantum mechanics to make predictions on a certain chemical system, a solution should be found to the time-independent Schrödinger equation:

$$\hat{H}\psi = E\psi \quad (1)$$

with E being the electronic energy, $\psi = \psi(x_1, x_2, \dots, x_n)$ the wavefunction, and \hat{H} the Hamiltonian operator:

$$\hat{H} = \sum_{i=1}^N \left(-\frac{1}{2} \nabla_i^2\right) + \sum_{i=1}^N v(\mathbf{r}_i) + \sum_{i < j}^N \frac{1}{r_{ij}} \quad (2)$$

with the “external” potential v acting on electron i due to nuclear charges Z_α being:

$$v(\mathbf{r}_i) = -\sum_{\alpha} \frac{Z_{\alpha}}{r_{i\alpha}} \quad (3)$$

The Hamiltonian basically consists of three factors: 1) a kinetic energy term T , 2) a potential energy term describing the electron-nucleus attraction V_{ne} and 3) a second potential energy term for electron-electron repulsion V_{ee} :

$$\hat{H} = T + V_{ne} + V_{ee} \quad (4)$$

There is currently no method in existence that can exactly solve the Schrödinger equation. Instead, the variational principle is used to approach the solution systematically. The variational principle states that the energy that is calculated using trial wavefunctions ψ_{trial} results in an upper bound of the true energy, such that

$$\langle \psi_{\text{trial}} | \hat{H} | \psi_{\text{trial}} \rangle = E_{\text{trial}} \geq E_0 = \langle \psi_0 | \hat{H} | \psi_0 \rangle \quad (5)$$

and E_{trial} can only equal E_0 if ψ_{trial} exactly matches ψ_0 . This means that function ψ_0 will give the lowest energy, E_0 . This approach, however, would require searching through all possible N -electron wavefunctions for that system.

The Hartree-Fock method approximates the N -electron wavefunction by taking the antisymmetrized product of N one-electron wavefunctions $\chi_i(x_i)$, also called a Slater determinant, Φ_{SD} , given here as the short-hand notation:

$$\Phi_{SD} = \frac{1}{\sqrt{N!}} \det\{\chi_1(r_1)\chi_2(r_2)\dots\chi_N(r_N)\} \quad (6)$$

The one-electron wavefunctions are spin orbitals, consisting of a spatial orbital term $\varphi_i(\mathbf{r})$ and of spin functions $\alpha(s)$ or $\beta(s)$, which are orthonormal. The variational principle can then be used to find the Slater determinant that gives the lowest energy. The presence of the Fock operator which depends on the spin orbitals/the eigenvalue problem solutions that need to be solved, gives rise to an eigenvalue equation that can only be solved iteratively through a self-consistent field (SCF) process. As mentioned before, due to the fact that the Slater determinant cannot exactly match the ground

state wave function, the calculated energy is always larger than the true ground state energy E_0 . The difference between the Hartree-Fock derived energy and the true ground state energy is called the correlation energy E_{HFC} . Contributing to this correlation energy is the dynamical electron correlation, caused by electron–electron repulsion. This is calculated to be too large by the Hartree-Fock procedure, because only the average electrostatic interaction is taken into account. This causes the electrons to move too close to one another. Secondly, there is a non-dynamical correlation, which stems from the Slater determinant not approximating the true ground state, which should ideally be expressed as a linear combination of many different determinants that represent different states. Many higher level methods have been developed in the past to correct these shortcomings that are in principle well-understood.

These post-Hartree-Fock methods can produce highly accurate results, but are practically not useful for models of realistic size, as the computational cost becomes intractable quickly. Whereas the simple Hartree-Fock method scales as N^3 , where N is the number of electrons, post-Hartree-Fock methods can scale as high as N^6 . By using a quantity such as the electron density, as opposed to wavefunctions, the amount of variables can be reduced to simply $3N$, because a density only depends on the 3D space. Furthermore, the electron density is an observable quantity, through methods such as X-Ray Diffraction, which wave functions are not. It can be shown that the electron density intrinsically contains all the key quantities needed to create the Hamiltonian.

Although there are earlier theories that utilize the electron density to calculate the ground state energy, such as the Thomas-Fermi model, the Hohenberg-Kohn theorem proves that the electron density indeed gives rise to a unique functional for $V_{\text{ext}}(\mathbf{r})$ and therefore the Hamiltonian \hat{H} . The electron density is given in terms of ψ as:

$$\rho(r) = N \iint |\psi(x_1, x_2, \dots, x_N)|^2 dx_1 dx_2 \dots dx_N \quad (7)$$

The energy can then be rewritten from equation (2.7) as:

$$E_0[\rho_0] = \int \rho_0(r) V_{ne} dr + T[\rho_0] + E_{ee}[\rho_0] \quad (8)$$

in which the first term is system dependent and the last two terms are system independent and form the Hohenberg-Kohn functional $F_{\text{HK}}[\rho_0]$. This Hohenberg-Kohn functional gives the expectation value $\langle \Psi | T + V_{ee} | \Psi \rangle$, the sum of the kinetic energy and electron–electron repulsion. This equation can still never lead to the exact energy, as the exact form of both the kinetic energy $T[\rho]$ and the electron-electron interaction, $E_{ee}[\rho]$, are not explicitly known.

Hohenberg and Kohn postulated a second theorem equivalent to the variational principle, stating that the calculated energy can only be the ground state energy if the density used is the ground state density, ρ_0 , expressed as:

$$E_0 \leq E[\rho_{trial}] + E_{ne}[\rho_{trial0}] + E_{ee}[\rho_{trial}] \quad (9)$$

Finally, the ground state energy, E_0 can be expressed as:

$$E_0 = \min_{\rho \rightarrow N} (F[\rho] + \int \rho(r) V_{Ne} dr) \quad (10)$$

where $F[\rho]$ represents the universal formula $E[\rho] = \min_{\psi \rightarrow N} \langle \psi | T + V_{ee} | \psi \rangle$, which is defined for all densities originating from antisymmetric wavefunction Ψ and contains the kinetic energy $T[\rho]$, classical Coulomb interaction $J[\rho]$, and non-classical interaction $E_{ncl}[\rho]$, such as self-interaction. Only the Coulomb interaction is explicitly known. Previous attempts had been made at defining the kinetic energy explicitly, but the results of these proved to be unstable and inaccurate. Kohn and Sham set out to solve this problem.

Kohn and Sham implemented a drastic simplification of the model by ignoring electron-electron interactions. This is expressed in Kohn-Sham orbitals as:

$$\rho(r) = 2 \sum_i^{occ} |\psi_i(r)|^2 = \rho_0(r) \quad (11)$$

On the other hand, by using this non-interacting system, Kohn and Sham realized that the true kinetic energy of the true system could not be calculated exactly. Rather, they used a formula that could be used to calculate the kinetic energy of the non-interacting system only:

$$T_s = -\frac{1}{2} \sum_1^N \langle \varphi_i | \nabla^2 | \varphi_i \rangle \quad (12)$$

The difference between the true kinetic energy T and the kinetic energy of the non-interacting system T_s , can then be added to a new term called the exchange-correlation energy, E_{xc} , which also includes the non-classical electrostatic interactions, E_{ncl} . Were the exchange correlation energy to be explicitly known, the Kohn-Sham approach would lead to the exact ground state energy, as opposed to the Hartree-Fock approach. Current research into density functional theory, therefore, focuses largely on the quest for the best approximation of the exchange correlation energy.

This section was summarized from *A Chemist's Guide to Density Functional Theory* by Wolfram Koch and Max C. Holthausen.²

1. Dehydrogenation of Linear and Cyclic Alkanes on 2D layered TiS₂

1.1 Introduction

Linear and cyclic alkenes, especially α -olefins, are high value reactants in industry, the vast majority of which are produced via steam cracking (SC).^{3,4} SC is not a selective process, but rather produces a mixture of different hydrocarbons, making it difficult to precisely tune into market demands for specific chemicals.^{5,6} Furthermore, as governments and consumers are pushing industry to search for cleaner, more sustainable solutions, alternatives to the highly energy intensive SC process are becoming increasingly important. Recent decades have seen a rise in the development of homogeneous catalysts that are capable of dehydrogenating alkanes, typically employing transition metals such as iridium, rhodium and ruthenium.⁷⁻¹² These catalysts form hydrides which may induce olefin isomerization, decreasing the selectivity towards the desired α -olefins. In 2017, Mindiola and coworkers reported the use of a titanium complex in tandem with a carbene transfer reagent to afford terminal alkenes, as well as some cyclic alkenes, with very high selectivity, since the Ti-complex does not form any hydrides, while using mild conditions.¹³ Homogeneous catalysts, however, have some drawbacks, such as the need for solvents in many cases, as well as laborious and energy intensive separation methods required to collect the product in high purity. Employing heterogeneous catalysts instead provides good solutions to a number of these problems.

Metal oxides have been used to dehydrogenate alkanes such as propane and ethane. Cr₂O₃/Al₂O₃ and Pt-Sn/Al₂O₃ are widely used in industry for alkene dehydrogenation.¹⁴ These metals, however, are either expensive or have toxic qualities, making them less favorable than other earth-abundant and less toxic alternatives. Furthermore, metal oxides tend to overoxidize alkanes, as the added value products in these reactions are more reactive under the reaction conditions. This increases selectivity to CO_x-type products.¹⁵⁻¹⁷ One solution to this problem is to use sulfur as a milder oxidant than oxygen.^{18,19}

Metal sulfides belong to a class of materials referred to as transition metal chalcogenides (TMCs), usually denoted MX₂ (X = S, Se or Te). TMCs have a variety of applications, such as electro- and photocatalysis and energy storage.²⁰⁻²² TiS₂, amongst others, was found to activate the C-H bond in methane and oxidatively couple carbon to form a new C-C bond in ethylene.²³ Activity and selectivity of these TMC catalysts was determined by the M-S bond strength; materials with stronger M-S bonds exhibited lower -CH₂ coupling activation barriers, while the C-H activation barrier

became higher. Unfortunately, this reaction requires temperatures over 1200 K to achieve sufficient yields and selectivities. Since titanium was also employed as a catalyst for alkane dehydrogenation under mild conditions, however, improvement of the TiS₂ nanosheet could lead to better activity and selectivity.

As for most TMCs, catalytic activity on TiS₂ occurs on the edge sites, and not on the basal plane. This could limit turnover if the edges site/basal plane ratio is very low.^{24–26} Cheon and coworkers found that several TMCs, including TiS₂, can be activated in the basal plane by reaction with Lewis acids such as AlCl₃.²⁷ This could improve the activity of the material as a catalyst. Although Marks and coworkers only probed C–H activation of methane, other literature suggests that larger alkanes, such as propane and isobutane may also be dehydrogenated by TiS₂.²⁸

In this study, the potential mechanism of the dehydrogenation of linear and cyclic alkanes was investigated using Density Functional Theory, in order to elucidate the viability of performing this reaction experimentally. The dehydrogenation mechanism was studied on two representative reactants: cyclohexane, representing cyclic alkanes, and *n*-butane, representing linear alkanes. Since there are, to our knowledge, no previous studies on the dehydrogenation of cyclic or linear alkanes on TMDs, this investigation may provide novel insights aiding future investigations into experimentally performing this reaction.

1.2 Theory

1.2.1 Transition Metal Dichalcogenides

Transition Metal Dichalcogenides (TMDs), denoted as MX₂, where X = S, Se or Te, are a class of material consisting of a sandwich-like structure in which the transition metal (TM) is sandwiched between two chalcogen layers. Usually these “sandwich” layers are held together by Van der Waals forces in the bulk material.^{29,30} This layered structure also allows for various methods to synthesize TMD nanosheets consisting of several or only one layer, such as exfoliation or vapor deposition.^{30–32} TMDs have a variety of applications, including in catalysis, semiconductors and sensors.^{20,33,34}

TMDs exist in several naturally occurring phases, the most common being octahedral (1T) and hexagonal (1H). Through phase engineering, alternative, non-thermodynamic phases, also called meta-stable phases, can be accessed, which can give rise to a change in the properties of the TMD.³⁵ The thermodynamically most stable phase of TiS₂ is the octahedral (1T) phase, referred to as 1T-TiS₂ (Fig. 1.2.1).

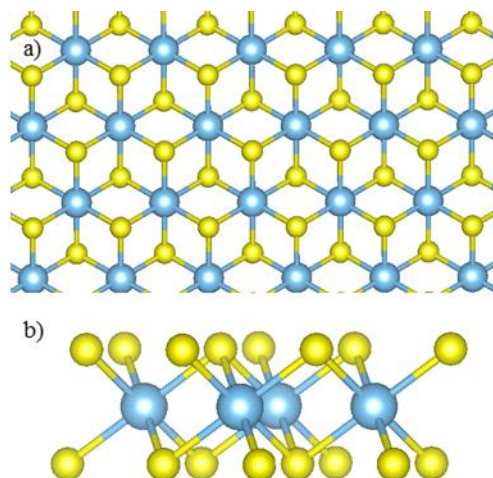


Fig. 1.2.1. The naturally occurring phase of TiS₂, octahedral or 1T-TiS₂, from a) top view onto the basal plane, and b) side view. Titanium = light blue; sulfur = yellow.

1.2.2 Model

Previous studies on MoS₂ nanosheets show that the shape of the nanosheet depends on the edge structure, which in turn is influenced by how sulfur-rich the environment is. In general, the most common edge structures are the zigzag edge and the armchair edge, which also occur for TiS₂. The zigzag edge structure is characterized by being S-rich, although S-poor zigzag edges, for example edges where only 50% sulfidation or no sulfidation at all, are also found. For MoS₂, the shape of the nanoflake can be influenced/changed by varying the chemical potential from Mo-rich to S-rich, which results in the mixture of zigzag/armchair edges to change to only zigzag edges, changes the nanosheet shape from dodecahedron to hexagon to triangle.

Since reactions of various TMDs with Lewis acids occur on the basal plane of the nanosheet, the edges of these TMDs are likely also Lewis acidic. Furthermore, both Zhu et al. and Li et al. calculated the stability of the occurring TiS₂ edge structures under varying chemical potential.^{23,36} Both investigations found that, except for very high chemical potential, the 50% sulfided zigzag edge state, ZZ1, is the most stable, as can also be seen in Fig. 1.2.2. For this reason, the calculations on the dehydrogenation mechanism were performed on a TiS₂ cluster with exclusively the ZZ1 edge structure. Other edge structures, such as the 100% sulfided edge (ZZ2), the 0% sulfided edge (ZZ0), and the armchair edge (AC), were tested to compare the adsorption energies of the alkanes. Potentially, changing the edge structure could also change the reactivity of the TMD nanosheet.

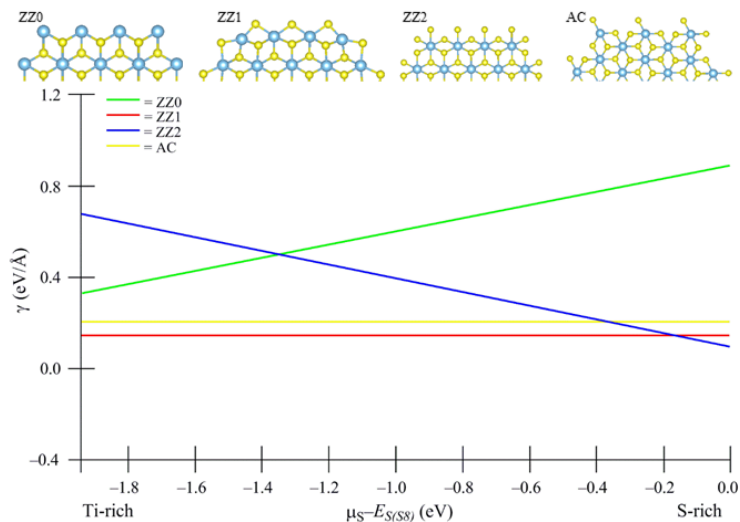


Fig. 1.2.2. The stability of the various 1T-TiS₂ edge structures, ZZ0, ZZ1, ZZ2 and AC. The figure is adapted from Li et al.;³⁶ similar results were also obtained by Zhu et al.²³ The stability was calculated using $\mu_{Ti} + 2\mu_S = E_{TiS_2(NS)}$, with μ_S being determined by the S₈ molecule for a S-rich environment, and the bulk Ti phase for a Ti-rich environment.

Zhu et al. mention that the C–H activation, as well as subsequent steps, occur on an S–S couple on the TiS₂ edge,²³ however this is unlikely to be the result of atomic S₂ adsorption on the edge, due to the above mentioned stability issue. Rather atomic sulfur is required to replenish the S-atoms on the nanosheet edge, as the reaction produces H₂S as a product also.

1.2.3 C–H Activation

C–H bond activation is an important chemical process, as it allows for the transformation of alkanes into value-added products in a very direct manner. Alkanes, of course, are widely available, but they are very inert, making C–H activation difficult to achieve. Mostly, production of value-added compounds, such as alcohols and carboxylic acids, uses alkenes as reactants, as they are much more reactive than alkanes. As mentioned, in the introduction, however, alkenes are in turn produced by the highly energy intensive steam cracking process. Therefore, accessing these value-added compounds directly through the naturally occurring alkanes, could be much more efficient and sustainable.

In general, three major pathways to activate C–H bonds can be identified: 1) homolytic cleavage or a radical pathway:



2) heterolytic cleavage, whereby a proton is abstracted from the alkane:



and 3) heterolytic cleavage with hydride abstraction from the alkane:



Here, these three mechanisms will be described briefly. For a more in depth overview of the various alkane C-H activation pathways, the reader is referred to a review on alkane CH activation pathways with organometallic compounds.³⁷

Homolytic Cleavage

C-H bonds undergo homolytic cleavage under attack from a, usually organic, radical. The radicals contain a SOMO, or singly occupied molecular orbital, which can either attack the HOMO or the LUMO of a bond. In the case of a C-H bond, the σ and σ^* orbitals are available, but due to the high energy of the σ^* orbitals, the SOMO will interact with the σ orbitals, or the HOMO. The radical in this case, acts as an electrophile. One example of this is the oxygen radical often used in biological processes, such as a hydroxyl or alkoxy radical.³⁸

A well-known example of homolytic C-H bond cleavage is the rebound mechanism.³⁹[ref] The hydrogen is abstracted from the reactant by the oxygen moiety, leaving the alkyl radical. This radical can then attack the newly formed hydroxide moiety, generating an alcohol complex. The rebound mechanism is commonly found in nature with the cytochrome P450 enzyme.^{38,40-42}

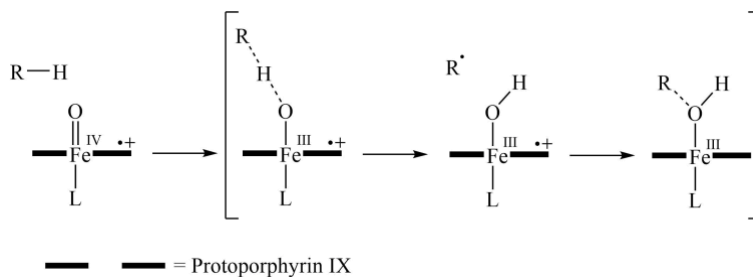


Fig. 1.2.3. The rebound mechanism as a pathway to C-H activation. Adapted from ref. 41.

Heterolytic Cleavage of Proton

Amongst the two heterolytic cleavage pathways, the pathway following proton abstraction is probably more common than that with hydride abstraction. A simple example of heterolytic C-H bond cleavage with proton abstraction is σ -bond metathesis.^{43,44} This heterolytic cleavage is often catalyzed by organometallic, M-X type complexes, where X is a heteroatom. This results in M-C and X-H type products.

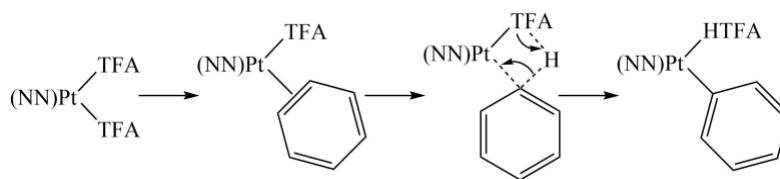


Fig. 1.2.4. The σ -bond metathesis mechanism as a pathway to heterolytic cleavage via proton abstraction. Adapted from ref. 44.

Heterolytic Cleavage of Hydride

A second heterolytic cleavage pathway of C–H bonds abstracts not a proton, but a hydride. This pathway is usually found for organometallic complexes with extremely electron-poor metal centers.⁴⁵ An example of heterolytic cleavage of the C–H bond via hydride abstraction in β -hydride elimination. This usually leads to a metal-hydride, but the hydride can also be concertedly transferred to another substrate shown below in Fig. 1.2.5.⁴⁶

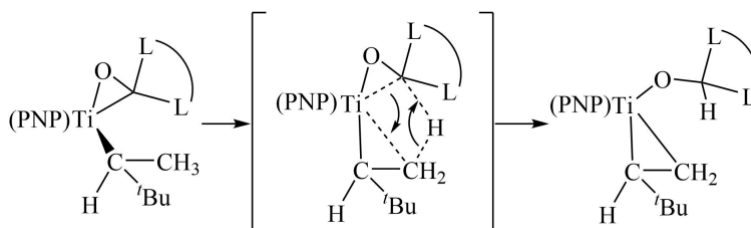


Fig. 1.2.5. The concerted β -hydride elimination-hydride transfer mechanism as a pathway to heterolytic cleavage via hydride abstraction. Adapted from ref. 46.

1.2.4 Proposed Mechanism

The reaction mechanism presented here is based largely on the reaction mechanism explored in the Zhu et al. paper of 2013. Cyclohexane and *n*-butane are used as representative substrates. The proposed mechanism for cyclohexane dehydrogenation is presented in Fig. 1.2.6. The full catalytic cycle is shown in Fig. 1.2.6a), including the expected intermediates and transition states. As shown in Fig. 1.2.6b), there are two basic steps: 1) the first C–H activation step over two S-atoms, and 2) the second C–H activation step of the adsorbed alkyl on another S-atom, which eliminates the alkene and H₂S.

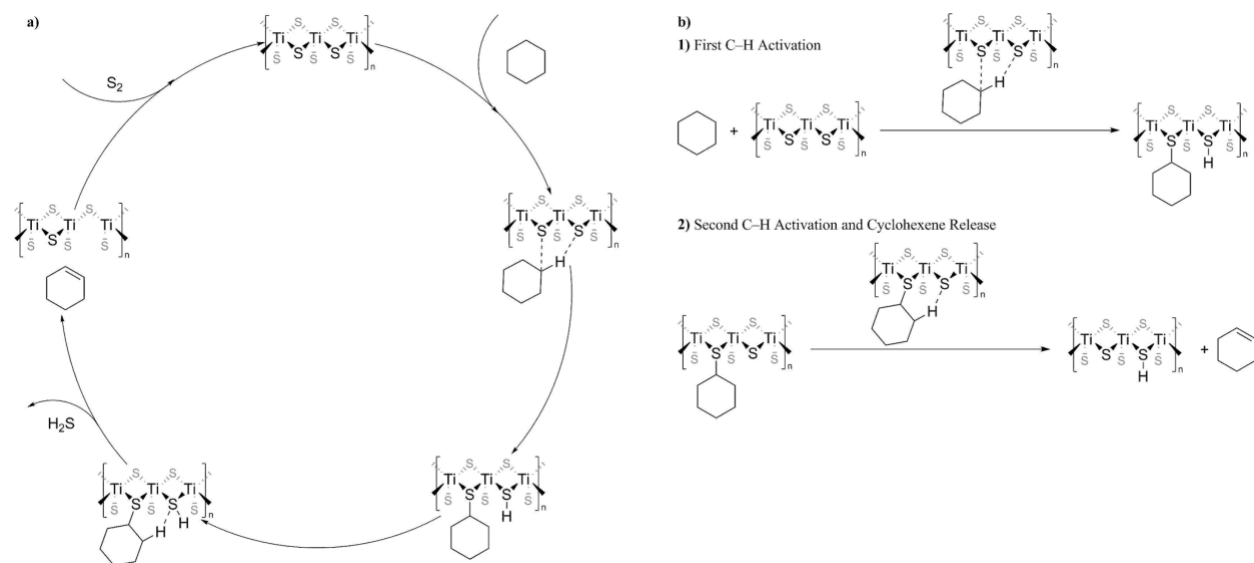


Fig. 1.2.6. a) The proposed catalytic cycle of cyclohexane dehydrogenation to cyclohexene, b) the individual C–H activation steps shown as they were applied in the calculation. The full catalytic cycle, including H₂S production, was not considered. Instead the C–H activation steps were investigated separately. The ZZ1 edge structure of the 1T-TiS₂ nanosheet is depicted schematically, with only the nearest neighbors of the edge-based Ti-atoms showing for clarity.

Cyclohexane is a relatively straightforward substrate; due to its symmetry and the fact that each carbon atom in the cycle has an identical environment, there is only one possible product. Butane dehydrogenation on the other hand, can afford both 1-butene and 2-butene, even when isomerization is not taken into account. Therefore, there are two possible pathways through which to achieve the first C–H activation step, shown in Fig. 1.2.7 a) and b), and three pathways for the second C–H activation, as shown in Fig. 1.2.7 c), d) and e).

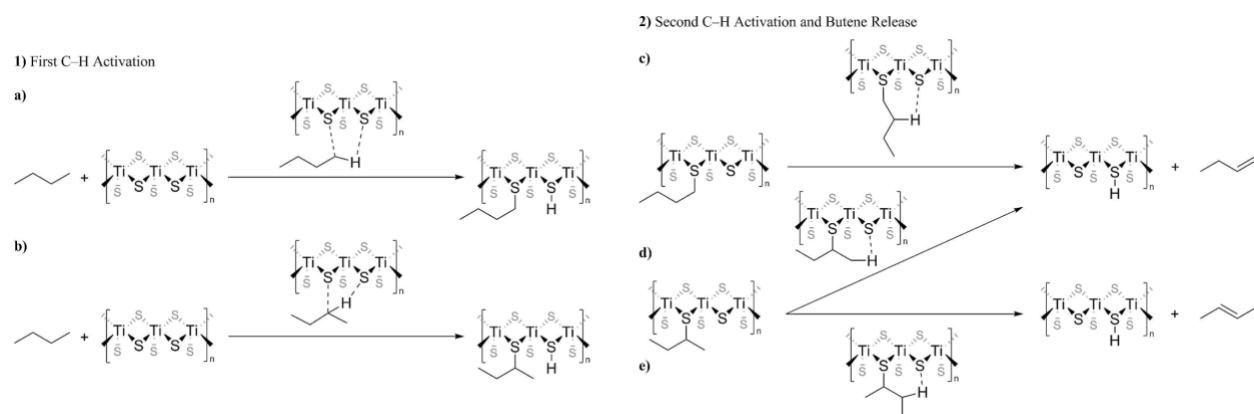


Fig. 1.2.7. The possible pathways for butane dehydrogenation, affording 1-butene or 2-butene. The first C–H activation can result in adsorption on a) the C1 position, or b) the C2 position. Subsequently, the second step can lead to the production of 1-butene via C–H activation on c) the C2 position or d) the C1 position, whilst C–H activation on e) the C3 position will lead to 2-butene.

As indicated, both of these C–H activations are assumed to occur over two different S-atoms, based on the investigations by Marks and coworkers, and Shan and coworkers.^{23,28} This means that a homolytic cleavage or radical mechanism for C–H activation was considered.

1.3 Computational Details

The intermediates and transition states were all located with the density functional based tight binding (DFTB) program as implemented in the Amsterdam Modeling Suite (AMS).⁹³ The GFN1-xTB functional was used for geometry optimization and single point energy calculation.⁴⁷ Recently, Grimme and coworkers have published a second version of the extended tight-binding functional, GFN2-xTB, which is more accurate than the first version;⁴⁸ however, this new functional was not available in AMS at the time of this investigation. Extended tight-binding functionals are advantageous over the more common self-consistent charge density functional tight-binding (SCC-DFTB) in that they do not use precalculated integrals, requiring the appropriate parameter files. Instead, extended tight-binding uses an extended Hückel-like Hamiltonian approximation, together with Slater-type orbitals. An energy cutoff of $1.0 \cdot 10^{-5}$ was employed. No periodic boundary conditions were applied to the 1T-TiS₂ hexagonal nanosheet cluster. The orbital occupation strategy was set to automatic.

Attempts were made at locating the related transition states by performing a potential energy surface (PES) scan along the reaction coordinate of the expected transition states, and subsequently running a transition state search from an initial guess taken from the PES scan. Unfortunately, this method proved to be unsuccessful, as transition states could not be located from any initial guess. Therefore, transition states were not considered in the discussion.

1.4 Results & Discussion

In this section the dehydrogenation mechanism of the two representative substrates, cyclohexane and *n*-butane, will be presented and discussed. First, the cyclohexane dehydrogenation mechanism will be examined, as it is the more straightforward of the two. Then, the *n*-butane dehydrogenation mechanism will be discussed. Transition states were not located, instead literature and other chemical precedents were used to rationalize which mechanism is most favorable.

1.4.1 Cyclohexane Dehydrogenation Mechanism

Fig. 1.4.1 shows the reaction energy profile of cyclohexane dehydrogenation calculated by DFTB as implemented in AMS. Cyclohexane initially coordinates to the TiS₂ sheet via interaction of a H-atom to an edge-based Ti-atom in **B**, resulting in an elongation of the C–H bond from 1.095 Å to 1.132 Å, shown in Fig. 1.4.2. This coincides with an interaction energy of –10.6 kcal/mol. Further C–H activation yields adduct **C**, which is more stable than **B** by 3.6 kcal/mol, at –14.2 kcal/mol. This is in stark contrast with the findings by Marks and coworkers, who found that the initial C–H activation of methane corresponds to a reaction energy of +3.6 kcal/mol. This addition was also calculated to have an activation barrier of 61.9 kcal/mol.²³ The C–H bond of methane is known to be the strongest bond amongst the hydrocarbons. This could explain why the addition of cyclohexane to the TiS₂ nanosheet leads to a more stable product. Furthermore, although a transition state was not located in this investigation, it is likely that the C–H activation barrier of cyclohexane is lower than that of methane.

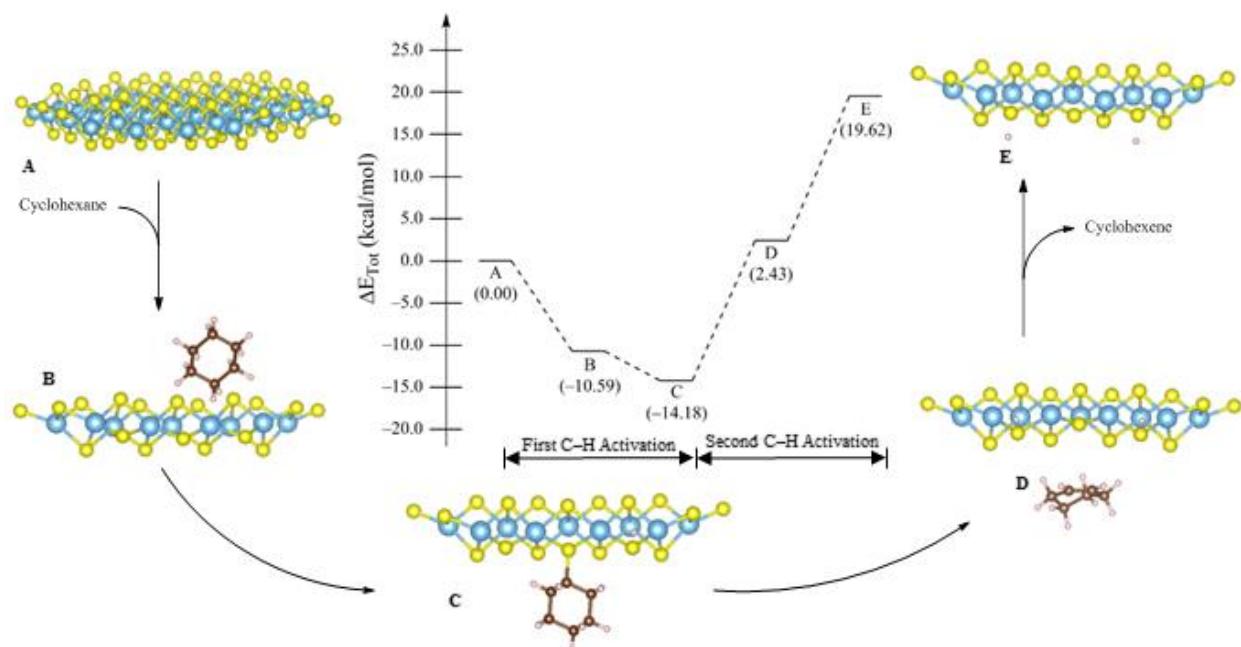


Fig. 1.4.1. The reaction energy profile of cyclohexane dehydrogenation on a 1T-TiS₂ hexagonal nanosheet with ZZ1 edge. The full nanosheet is shown in **A**, but for the other geometries only the first two rows of titanium on the edge where the reaction took place, are shown for clarity. Titanium = light blue; sulfur = yellow; brown = carbon; white = hydrogen.

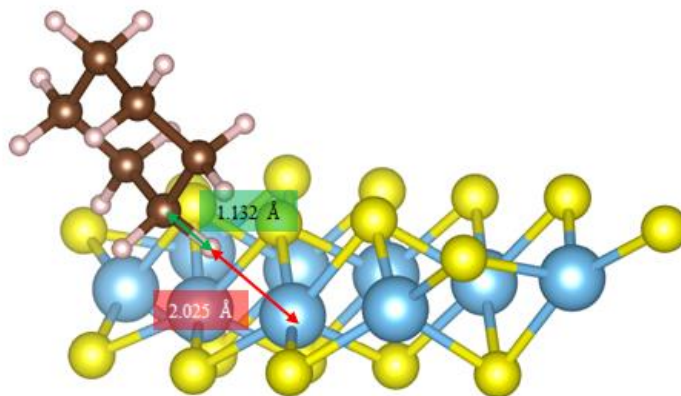


Fig. 1.4.2. C–H bond length and H···Ti distance after coordination of cyclohexane to the 1T–TiS₂ nanosheet in **B**, with $\Delta E_{\text{Tot}} = -10.59$ kcal/mol. Titanium = light blue; sulfur = yellow; brown = carbon; white = hydrogen.

The second dehydrogenation step was found to be comparatively unfavorable. The thiirane intermediate proposed by Marks and coworkers²³ was found to be highly unfavorable for cyclohexene, shown in Fig. S1.1, at 41.2 kcal/mol. Instead direct release of cyclohexene is likely more favorable for this compound. Therefore, intermediate **D**, with dissociated cyclohexene, was located. The formation of this intermediate is associated with an increase in energy of 16.6 kcal/mol, increasing the energy of **D** to $\Delta E_{\text{Tot}} = 2.4$ kcal/mol. Interestingly, the C=C bond in **D** is 1.327 Å, which is shorter than the normal cyclohexane C=C bond at 1.366 Å, by approximately 0.03 Å. While this is not a dramatic change in bond length, it is noteworthy. The reason for the C=C bond shortening remains, as of yet, unknown.

The reaction energy of the final step, cyclohexene elimination affording **E**, was 17.2 kcal/mol, increasing the ΔE_{Tot} to 19.62 kcal/mol. It is well established that cyclohexene is comparatively less stable than cyclohexane. For example, during benzene hydrogenation almost only the fully hydrogenated cyclohexane is observed as the product, and not the partly hydrogenated cyclohexadiene or cyclohexene.⁴⁹ Similar observations are made for cyclohexane dehydrogenation on other catalysts.⁵⁰ Currently, cyclohexene is commercially produced by Asahi with a Ru–Zn/SiO₂ catalyst from benzene hydrogenation;⁵¹ however, this process suffers both from poor yield and poor selectivity, indicating the difficulties of cyclohexene synthesis. It should, therefore, not be surprising that formation of cyclohexene is energetically unfavorable.

In the second C–H activation step the hydrogen was assumed to be adsorbed on another edge-based S-atom than in the first C–H activation step. Conceivably, the second C–H activation step may occur on the same S-atom, directly followed by H₂S elimination. H₂S production was observed in several experimental investigations on alkane C–H activation on TMDs,^{23,28} and may act as a driving

force for the reaction, as H₂S gas production would be entropically favorable. This mechanism was also considered and is shown in the supporting information (Fig. S1.2). Contrary to expectations, however, the release of H₂S results in a sharp increase in the reaction energy of 46.3 kcal/mol. Comparison of the TiS₂ nanosheet before and after H₂S elimination reveals that the TiS₂ nanosheet is strongly destabilized by the removal of a S-atom from the edge. The final product **E**, neglecting H₂S formation, is more stable than final product **G'**, including H₂S formation, by 46.3 kcal/mol. The calculated 1T-TiS₂ edge structure stability already hinted at this, as the ZZ0 edge, containing bare Ti-edge sites, is one of the least stable edge structures, as shown in Fig. 1.2.2.^{23,36} This result does not necessarily question the H₂S formation observed in experiments; there could simply be a different mechanism in place leading to H₂S production. Furthermore, addition of elemental sulfur, as mentioned in ref. 23, could drive the release of H₂S from the 1T-TiS₂ edge. Further investigations would be necessary to determine this; however, this was outside of the scope of this research.

1.4.2 Butane Dehydrogenation Mechanism

The reaction energy profile of *n*-butane dehydrogenation is presented in Fig. 1.4.3., showing both the formation of 1-butene and 2-butene. The butane initially coordinates to the edge-based Ti-atom of the TiS₂ nanosheet edge through the C2-H bond in **2**, elongating the bond from 1.095 Å to 1.134 Å, as shown in Fig. 1.4.4. Then, either C1-H activation or C2-H activation can take place to form the butane adducts **3** and **3'**, respectively. Considering that butane interacts with the TiS₂ nanosheet via the C2-H bond, this activation may be easier than C1-H activation. Nevertheless, the reaction energy for **3** is -12.7 kcal/mol, while that for **3'** is -12.0 kcal/mol. This indicates that there is no significant thermodynamic preference for any adduct. Instead, the activation barrier for formation of **3'** may be lower than the activation barrier for **3**, due to the coordination of butane to the nanosheet via the C2-H bond. In order to conclusively determine this, the related transition states should be located and compared. Unfortunately, that was not possible during this investigation.

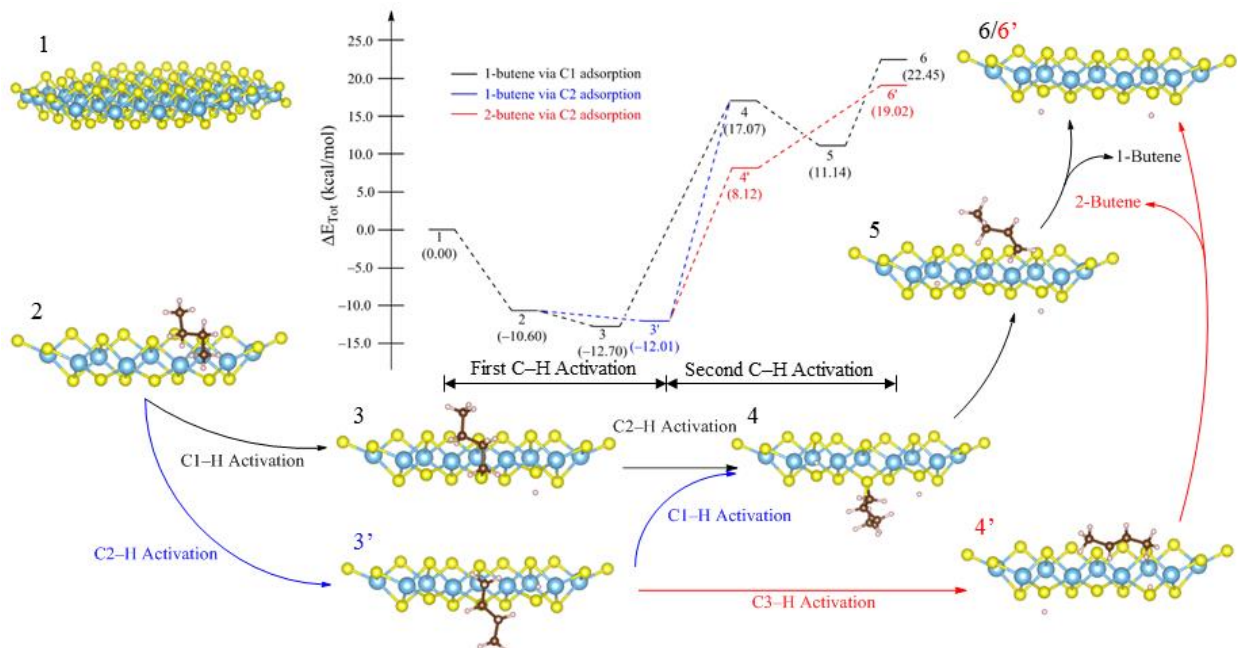


Fig. 1.4.3. The reaction energy profile of butane dehydrogenation on the 1T-TiS₂ nanosheet with ZZ1 edge. The full nanosheet is shown in **1**, while only the first two titanium rows on the edge where the reaction takes place, are shown in the other structures for clarity. Titanium = light blue; sulfur = yellow; brown = carbon; white = hydrogen.

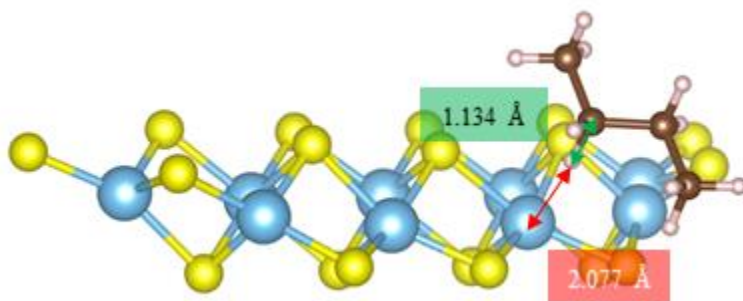


Fig. 1.4.4. C-H bond length and H...Ti distance after coordination of *n*-butane to the 1T-TiS₂ nanosheet in **2**, with $\Delta E_{\text{Tot}} = -10.60$ kcal/mol. Titanium = light blue; sulfur = yellow; brown = carbon; white = hydrogen.

The second dehydrogenation step from **3** occurs via C2-H activation, affording thiirane intermediate **4** with an increase in energy of 29.8 kcal/mol. Intermediate **4** can also be accessed by **3'** via C1-H activation, coinciding with an energy increase of 29.1 kcal/mol. This is an unexpectedly high energy difference, likely also coinciding with a high energy barrier. The thiirane intermediate **4** is a sterically highly strained conformation, likely contributing to this step energy increase. A more stable intermediate can be formed by C3-H activation from **3'**, affording **4'** with corresponding increase in energy of 20.1 kcal/mol. This intermediate affords 2-butene directly, without the formation of a thiirane intermediate, explaining the relative stability. An attempt at locating the corresponding thiirane intermediate proved unsuccessful, the resulting geometry of which is shown

in Fig. S1.3, and the intermediate was not considered in the final mechanism. The reason why this thiirane is not formed remains unclear.

The formation of the thiirane intermediate **4** seems to be a crucial step in the second C–H activation, yet there are C–H activations occurring on other C-positions imaginable, that may be associated with both lower activation barriers, as well as lower intermediates. This could be, for example, C3–H activation and C4–H activation from **3**, resulting in the corresponding 2-methylthietane and tetrahydrothiophene intermediates shown in Fig. 1.4.5c) and d), respectively. Both of these structures are more stable than thiirane intermediate **4** shown in Fig. 1.4.5a) by –3.1 and –14.1 kcal/mol, respectively. These intermediates, however, would not lead to stable products compared to intermediate **4**. A cyclopropane or cyclobutane would be formed, which are highly strained compounds that are unlikely to be stable compared to 1-butene or 2-butene. Therefore, the formation of these products is not considered likely.

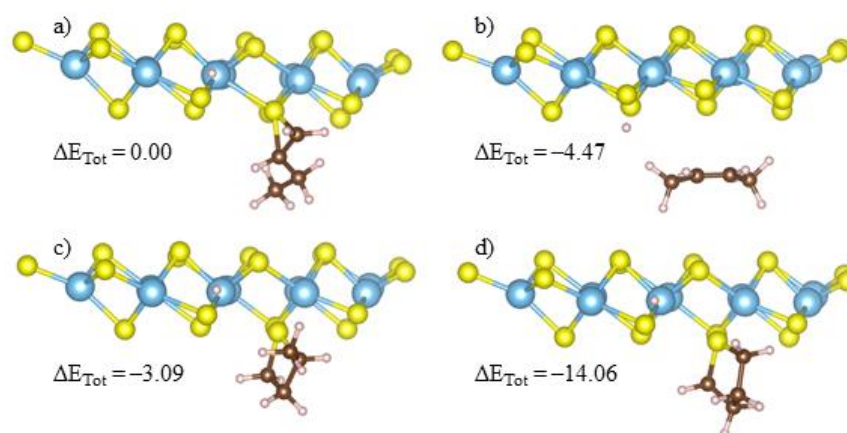


Fig. 1.4.5. The possible heterocycle intermediates afforded after C–H activation from **3**, where a) the expected thiirane intermediate **4** in 1-butene formation, b) the directly dissociated intermediate **4'** in 2-butene formation, c) a 2-methylthietane intermediate afforded after C3–H activation from **3** and d) a tetrahydrothiophene intermediate afforded after C4–H activation from **3**. ΔE_{Tot} in kcal/mol. Titanium = light blue; sulfur = yellow; brown = carbon; white = hydrogen.

Finally, 1-butene dissociates from the TiS_2 edge, allowing the energy to decrease by almost 6 kcal/mol. There appears to be an interaction between 1-butene and the TiS_2 nanosheet. Complete dissociation causes the energy to increase again to 22.5 kcal/mol in **6**, an increase of 11.3 kcal/mol relative to **5**, which is the interaction energy. This is also the case for **4'** and **6'**, which show an increase in energy from 8.1 kcal/mol to 19.0 kcal/mol, corresponding to an interaction energy of 10.9 kcal/mol. This calculated interaction energy is likely an artefact of the calculations being performed

under vacuum. The observed interaction between the 1T-TiS₂ sheet and the dissociated butene in **4'** and **5** may disappear when a solvent is added to the nanosheet during the calculation.

The mechanism including H₂S production was also analyzed for butane dehydrogenation and is shown in the appendix in Fig. S1.4. Again, the products of this pathway, **7** and **7'**, are significantly destabilized compared to the mechanism shown in Fig. 1.4.3, by 46.3 kcal/mol. This is exactly the same as the destabilization afforded to the cyclohexane dehydrogenation mechanism for H₂S production shown in Fig. S1.2, further strengthening the aforementioned suspicion that the TiS₂ nanosheet with one edge-based S-atom is responsible for the destabilization.

1.4.3. General Considerations

According to Marks and coworkers, C–H activation occurs over a S–S couple on the 1T-TiS₂ edge, whereby the hydrogen is adsorbed on one S-atom and the alkyl is adsorbed on the neighboring one.²³ This would be consistent with a radical pathway of C–H activation. It is also possible for C–H activation to occur on just one S-atom, in a similar fashion to the rebound mechanism.³⁹ These two modes of adsorption are shown in Fig. 1.4.6. Adsorption on the same S-atom has a higher energy than adsorption on two different S-atoms by 10.1 kcal/mol, clearly indicating that this adsorption mode is not feasible. One argument against the radical mechanism may be drawn from Fig. 1.4.2., where the C–H bond is shown to be elongated and pointing in the direction of an edge-based Ti-atom. This may indicate that titanium participates in C–H cleavage, perhaps via a heterolytic cleavage mechanism by abstracting a hydride. This mechanism is not further explored here, however, as it is not supported by literature.^{23,28} Furthermore, the observed H₂S production during experimental work, would suggest that the H-atoms are in fact adsorbed on the edge-based S-atoms.

It is not unimaginable that the edge-based S-atoms alone function to homolytically cleave the C–H bond. In TiS₂, as in TiO₂, the Ti-atoms are in an extremely electron-poor state, having an oxidation state of 4+ and no *d*-electrons. While the Ti(IV) state is very stable, the Ti-atoms are likely pulling electron density from the electron rich S₂-atoms. This could give the edge-based S-atoms some radical character. In order to verify if this is the case, single point calculations were performed in the Jaguar 9.1 suite of modeling programs to find the spin densities on the 1T-TiS₂ model nanosheet. Unfortunately, the calculations were not completed at the time of finishing this work.

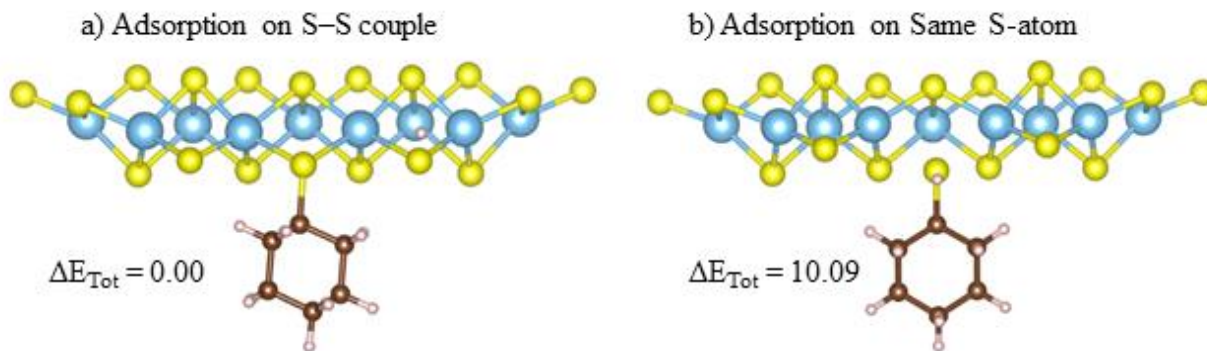


Fig. 1.4.6. Comparison of the products of C–H activation via a) the radical mechanism, and b) the rebound mechanism. ΔE_{Tot} is given in kcal/mol. Titanium = light blue; sulfur = yellow; brown = carbon; white = hydrogen.

The initial C–H activation step for both cyclohexane and *n*-butane dehydrogenation was found to be surprisingly favorable, compared to literature.²³ This may, however, simply be an artefact of the calculation. The calculations were performed in vacuum, rendering the TiS₂ nanosheet in a kind of “activated state”. During reaction in experimental conditions, the nanosheet would be placed in a solvent or surrounded by a gas. Solvation may give rise to thermodynamically favorable interactions, resulting in a stabilization of the solute, in this case the 1T-TiS₂ nanosheet. Therefore, while the current results show a reaction energy of –14.2 kcal/mol in vacuum for **C** formation, adding a solvent may stabilize the bare TiS₂ sheet to be more stable than **C**. Similarly, the observed interactions in **B** and **D** may be less significant, due to the TiS₂ surface being covered by solvent molecules. The same thing can be said for the dehydrogenation mechanism of *n*-butane.

There are methods to model a solvent: explicitly modeling the solvent by placing solvent molecules around the nanosheet, or implicitly by applying a continuum dielectric.⁵² Unfortunately, explicitly modeling the solvent is computationally expensive and requires boundary conditions, which were not applied in the calculations presented here. The implicit solvation model is vastly more efficient than the explicit model, however it is not available in DFTB.⁵³ Using another computational method, such as VASP or Jaguar 9.1, may provide a solution. This was, however, outside of the scope of this investigation. Nevertheless, it is important to keep these effects in mind when considering the implications of these results. While they may provide an initial understanding of the investigated system, and indicate the feasibility of achieving cyclohexane and *n*-butane dehydrogenation catalyzed by a 1T-TiS₂ nanosheet, a model system in vacuum will never fully approximate the experimental results.

The alkene formation in both mechanisms discussed above are accompanied by large reaction energies around 20 kcal/mol. Employing mild conditions in this reaction, as was suggested

to be possible in ref. 11, is unlikely to afford high yields in this system. In fact, Marks and coworkers report the reaction of methane C–H activation taking place at over 1200 K,²³ and Shan and coworkers heat the reaction set-up to more than 800 K to dehydrogenate isobutane over TMDs.²⁸ Under such reaction conditions, the above mechanism is much more likely to be achieved.

1.5 Conclusion & Outlook

In summary, the possibility of performing cyclohexane and *n*-butane dehydrogenation on a 1T-TiS₂ nanosheet were explored with Density Functional Tight-Binding by exploring the energetics of the potential reaction pathways. Initial C–H activation was found to be most favorable on an S–S couple on the TiS₂ edge. Subsequently, a second C–H activation takes place on the neighboring C-atom, directly affording the alkene in the case of cyclohexene and 2-butene, but resulting in a thiirane intermediate for the formation of 1-butane. Production of H₂S was found to be unfavorable due to the instability that is afforded to the 1T-TiS₂ nanosheet after removal of an edge-based S-atom. C–H activation on TiS₂ is thought to occur via the homolytic or radical pathway, but further investigation is required to confirm this. This can be investigated, for example, by calculating the spin densities on the edge-based S- and Ti-atoms. Since the calculations were performed in vacuum, solvation effects are not considered in this work. Future investigations may benefit from including solvent effects, to obtain energies that are more accurate to experimental measurements. To our knowledge there is no known experimental literature investigating the dehydrogenation of linear or cyclic alkanes in TiS₂; therefore, this work may provide a guideline to future attempts at performing this reaction experimentally or studying it computationally.

In this study, transition states were not located with DFTB. In order to gain a better understanding of the feasibility of these reactions, using a different method to locate transition states, such as the Nudged Elastic Band method, could provide useful insights. Furthermore, employing computational methods such as the Vienna Ab initio Simulation Package, which is known to be more accurate than DFTB, could improve the results obtained in this study.

Future investigations could focus on the activation of the basal plane of TiS₂, as mentioned in the introduction. Since layered TMDs such as TiS₂ are known to be active only on the edge sites, activating the basal plane could improve the activity of the nanosheet catalyst. Other modifications of the nanosheet, such as oxidation of the edges to form a TiO₂/TiS₂ are known to impact the electrochemical properties of the nanosheet. The impact of such modifications on the

dehydrogenation of linear and cyclic alkanes could also be of interest. Using these two modifications in tandem, may be able to improve not only the activity of the TiS₂ nanosheet, but also the selectivity.

2. The Electronic Structure of β -Diketimate Manganese Hydride Dimer

2.1 Introduction

Low-coordinate transition metal complexes have been extensively studied due to the reactivity enabled by their coordinatively and electronically unsaturated structures.⁵⁴ It has been widely acknowledged that bulky ligands can stabilize low-coordinate electron-deficient complexes by preventing aggregation.⁵⁵ One ubiquitous group of bulky ligands that support electron-deficient metal centers is the β -diketimate ligand, commonly abbreviated as BDI or nacnac. Practical advantages such as facile synthesis and modification⁵⁶ enable systematic studies on β -diketimate ligands.^{57,58} For instance, the Mindiola group reported catalytic carboamination with a titanium complex supported by a β -diketimate ligand.⁵⁹ The Holland group also reported a Fe(II) fluoride β -diketimate complex that catalyzes the hydrodefluorination of fluorocarbons.⁶⁰ Despite extensive investigation on transition metal complexes supported by β -diketimate ligands, (BDI)Mn complexes have only recently been found to exhibit catalytic activity. Notably, the bridging hydride dimer, $[(2,6\text{-}i\text{Pr}_2\text{PhBDI})\text{Mn}(\mu\text{-H})]_2$ (**1**), has been structurally characterized and found to catalyze the hydrosilylation of alkenes.¹

Intuitively, one can argue that complex **1**, which has 13 electrons per metal center, would prefer to engage in metal–metal bonding to compensate for its electron deficiency, as shown in Fig. 2.1.1. The existence of a metal–metal triple bond would formally provide three additional electrons to each Mn center, discounting expected weak δ -interactions, leading to a more reasonable electron count of 16. Few dimeric β -diketimate manganese complexes have been reported;^{61–63} however, a detailed analysis of their electronic structures has yet to be reported. In this contribution, the bond orders between the Mn centers of **1** and related complexes featuring bridging p-block elements are explored using a combination of computational and experimental techniques. Given that the coordinatively- and electronically-unsaturated nature of (BDI)Mn complexes has rendered them capable of catalyzing hydrofunctionalization reactions, a main goal of this research was to evaluate the electronic structure of **1** to disclose the origin of its catalytic activity.

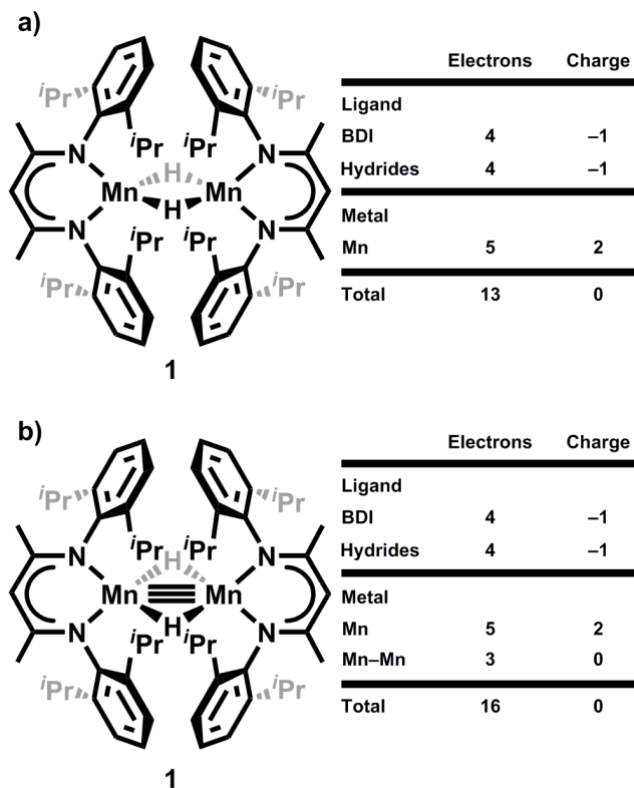


Fig. 2.1.1. Electron-count diagrams of a) **1** excluding metal–metal bonding and b) **1** including a triple bond between the manganese centers.

2.2 Computational Details

All calculations were performed using density functional theory (DFT) as implemented in the Jaguar 9.1 suite of ab initio quantum chemistry programs.⁶⁴ The B3LYP functional^{65,66} with Grimme's D3⁶⁷ dispersion correction (B3LYP-D3) was exploited as the standard, together with the 6-31G** basis set for geometry optimizations. The Los Alamos LACVP** basis including relativistic effective core potential was used to represent manganese.^{68–70} In order to obtain more reliable energies, single point calculations were performed on the optimized geometries using Dunning's correlation-consistent triple- ζ cc-pVTZ(-f) basis set for main group elements and LACV3P** for manganese.⁷¹ The zero-point energy (ZPE), entropic and solvation contributions to the Gibbs energy are obtained from the same level of theory as the geometry optimizations (B3LYP-D3/6-31G**/LACVP**). Energy decomposition was computed with the B3LYP-D3 functional using the Amsterdam Density Functional (ADF 2019) package.⁷² The optimized geometries were confirmed to be the local minima on the potential energy surfaces by showing the absence of an imaginary frequency. The solvation calculations utilized self-consistent reaction field (SCRF) approach on the gas phase geometry to model the solvation shell of dielectric constants $\epsilon = 2.284$ for benzene and $\epsilon = 2.397$ for toluene. As is

the case for all continuum models, the solvation energies are subject to empirical parametrization of the atomic radii that are used to generate the solute surface.⁷³ We employed the standard set of optimized radii in Jaguar for H (1.150 Å), C(1.900 Å), N(1.600 Å), O(1.600 Å), F(1.682 Å) and Mn(1.480 Å).

The energy components have been computed with the following protocol. The free energy in solution phase $G(\text{sol})$ has been calculated as follows:

$$G(\text{sol}) = G(\text{gas}) + G_{\text{solv}} \quad (2.1)$$

$$G(\text{gas}) = H(\text{gas}) - TS(\text{gas}) \quad (2.2)$$

$$H(\text{gas}) = E(\text{SCF}) + \text{ZPE} \quad (2.3)$$

$$\Delta E(\text{SCF}) = \Sigma E(\text{SCF}) \text{ for products} - \Sigma E(\text{SCF}) \text{ for reactants} \quad (2.4)$$

$$\Delta G(\text{sol}) = \Sigma G(\text{sol}) \text{ for products} - \Sigma G(\text{sol}) \text{ for reactants} \quad (2.5)$$

$G(\text{gas})$ is the gas phase free energy; $G(\text{solv})$ is the free energy of solvation; $H(\text{gas})$ is the gas phase enthalpy; T is the temperature (298.15 K); $S(\text{gas})$ is the gas phase entropy; $E(\text{SCF})$ is the electronic energy derived from the SCF method and ZPE is the zero-point energy.

In principle, multi-reference methods such as CASSCF are required to describe rigorously an antiferromagnetically coupled spin state of the dimer, which is impracticable for systems of this size because of computational demands. In practice, Noodleman's broken symmetry (BS) approach,⁷⁴ which makes use of the Heisenberg spin operator formalism to obtain a reasonable electronic structure description of transition metal dimers, provides a working protocol for single reference methods such as DFT employing the unrestricted spin formalism. We closely followed the protocol described elsewhere⁷⁵ to obtain the BS orbitals and used the unrestricted spin formalism in all calculations. Essentially, the valence bond descriptions of the molecules were used as initial guesses to generate a molecular wavefunction in terms of localized orbitals that undergo the SCF procedure. The large number of unpaired electrons at the manganese core and the intrinsic electronic flexibility of the transition metals give rise to a number of energetically similar electronic states. Therefore, we used spin densities derived from Mulliken population analysis that have also been employed in many previous studies⁷⁵ to confirm that the correct, or at least a reasonable, electronic state had been reproduced.

2.3 Results & Discussion

2.3.1 MO-Diagrams of Ferromagnetically and Antiferromagnetically coupled **1**

In order to gain an initial understanding of the electronic structure of **1**, conceptual MO diagrams were constructed (Fig. 2.3.1). Fig. 2.3.1a depicts the MO diagram of $\{[(2,6\text{-iPr}_2\text{PhBDI})\text{Mn}]_2\}^{2+}$, which lacks bridging hydride ligands. The in-phase d_{xy} , d_{xz} and $d_{x^2-y^2}$ orbital combinations could interact to form the expected Mn–Mn triple bond based on the electron count. The apparent non-bonding orbitals, the bonding and antibonding combinations of the d_{z^2} and d_{yz} orbitals, are expected to be nearly degenerate, giving rise to four unpaired electrons. This diagram forms the basis of the conceptual molecular orbital diagram of **1**, which is shown in Fig. 2.3.1b. For **1**, the in-phase combination of the d_{xy} orbitals is expected to interact with the out-of-phase combination of the hydride s-orbitals. After taking the bridging hydride ligands into account, four unpaired electrons remain in the weakly interacting d_{z^2} and d_{yz} orbitals. Importantly, the frontier molecular orbital diagram in Fig. 2.3.1b is consistent with the previously reported solution magnetic susceptibility of **1** ($5.2 \mu_B$ at 298 K).¹ The electrons initially expected to contribute to a bonding interaction between the metal centers, as shown in Fig. 2.3.1a, are instead redistributed into Mn–H bonds via a σ -bonding as illustrated in Fig. 2.3.1b. This σ -bonding interaction is speculated to reduce the bond order between the two manganese centers to be between 2 and 3. The crystal structure of **1**, however, reveals a Mn–Mn distance of 2.814 Å, which is significantly longer than the bond length expected for a Mn–Mn triple bond.^{76,77} Manganese dimers reported to feature Mn–Mn triple bonds are known to possess Mn–Mn distances of approximately 2.2 Å,^{78,79} suggesting that the bond order between the Mn centers of **1** is considerably lower than 3.

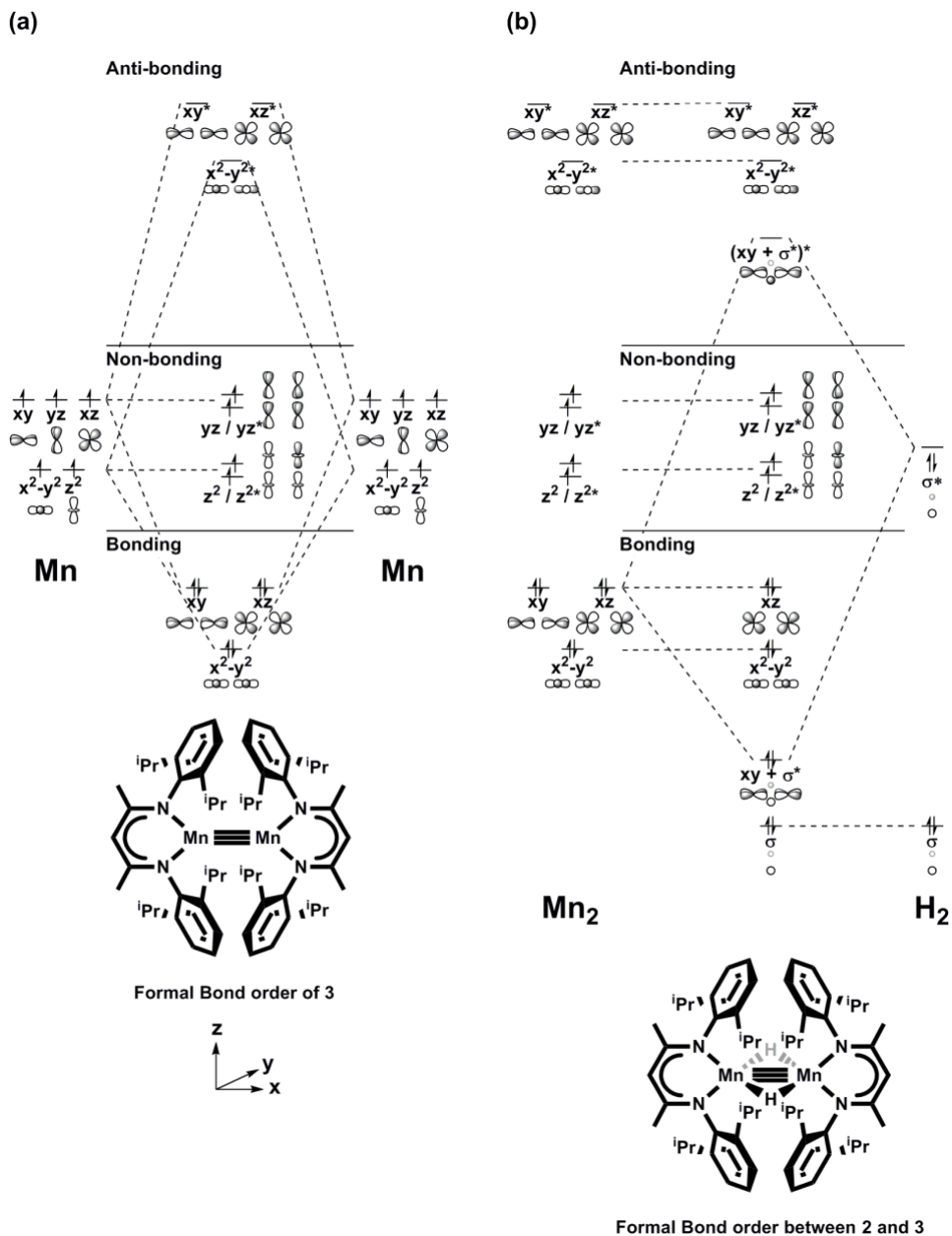


Fig. 2.3.1. Conceptual MO diagrams for a) $\{[(2,6\text{-}i\text{Pr}_2\text{PhBDI})\text{Mn}]_2\}^{2+}$ and b) **1**, which features an interaction between the out-of-phase hydride combination and the in-phase d_{xy} orbital combination.

It must be mentioned that the neutral analogue of $\{[(2,6\text{-}i\text{Pr}_2\text{PhBDI})\text{Mn}]_2\}^{2+}$, $[(2,6\text{-}i\text{Pr}_2\text{PhBDI})\text{Mn}]_2$, was reported by Roesky et al. in 2005 (Fig. 2.3.2).⁶² Notably, single crystal X-ray diffraction analysis revealed that $[(2,6\text{-}i\text{Pr}_2\text{PhBDI})\text{Mn}]_2$ features staggered BDI ligands and a Mn–Mn distance of 2.721 Å. This distance is shorter than the same distance in **1** (2.814 Å), but still considerably longer than what would be expected for a complex featuring a Mn–Mn triple bond. The contraction of the distance between metal centers was reported to attribute to a bonding interaction between 4s orbitals in metal centers. Instead of having a $3d_6$ configuration, the Mn(I) centers adopt a $4s^1 3d_5$ configuration;

therefore, two 4s orbitals from each metal center can form an in-phase combination. Further information about the ground state and magnetism of $[(2,6\text{-}i\text{Pr}_2\text{PhBDI})\text{Mn}]_2$ was gathered by DFT analysis;⁶² however, this compound's absence of Mn–Mn multiple bonding had not been investigated.

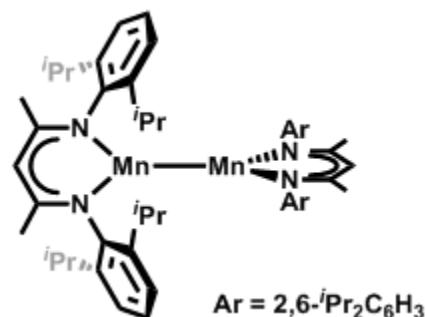
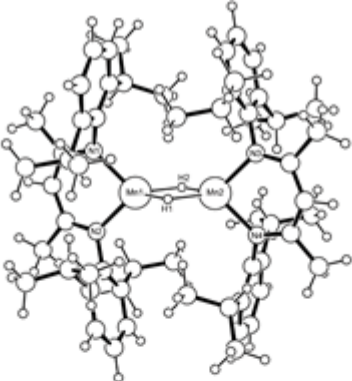


Fig. 2.3.2. The dimeric manganese complex, $[(2,6\text{-}i\text{Pr}_2\text{PhBDI})\text{Mn}]_2$, reported by Roesky et al.

The X-band (9.4 GHz) electron paramagnetic resonance (EPR) spectrum of **1** was reported in 2018.¹ At 106 K, a superposition of three spin states, $S = 1$, $S = 2$ and $S = 3$, was observed. The spectral analysis via computer simulation showed that the two Mn(II) centers are high spin and are antiferromagnetically coupled with each other. Likewise, we have found the $S = 0$ state to be the ground state and the electronic energy of the $S = 5$ state is higher than that of the $S = 0$ state by 0.9 kcal/mol, implying that the F-coupled state is accessible. Consequently, the spin states in between the two extreme states, $S = 1$, $S = 2$, $S = 3$ and $S = 4$ states, can be accessed as well at the EPR condition and a certain temperature. Considering the ambient temperature solution magnetic susceptibility data collected for **1** ($5.2 \mu_B$), the observed $S = 2$ spin state is also the ensemble average value of all possible states accessible at room temperature. Therefore, computational models with ferromagnetic (**1-F**) and antiferromagnetic coupling (**1-AF**) between the metal centers, representing the extreme states, were selected to evaluate.

The optimized geometries of **1-F** and **1-AF** are in good agreement with the crystal structure (**1-CS**), as shown in Table 2.3.1. The computed Mn–Mn distances are 2.871 Å for both **1-F** and **1-AF** cases, which is comparable to that of the crystal structure, 2.814 Å. The tetrahedral-like geometry around the Mn centers and the eclipsed positioning of the β -diketiminato ligands were clearly reproduced. Fig. 2.3.3 shows the structural consistency among the X-ray crystallographic data and the optimized geometries of **1-F** and **1-AF**. The absence of a multiple bond between the manganese centers was observed, with Mayer bond orders of 0.21 and 0.27 for **1-F** and **1-AF**, respectively.

Table 2.3.1. Optimized structural parameters of the dimanganese complexes **1-F**, **1-AF** and **1-CS**



	1-F	1-AF	1-CS
	calc.	calc.	exp.
Spin-Coupling	F	AF	-
Mn1-H1 (Å)	1.897	1.834	1.791
Mn1-H2 (Å)	1.919	1.923	1.938
Mn1-N1 (Å)	2.099	2.103	2.086
Mn1-N2 (Å)	2.116	2.110	2.098
∠H1-Mn1-N2 (°)	116.9	120.5	121.1
∠N2-Mn1-N1 (°)	90.1	90.5	91.5
∠N1-Mn1-H2 (°)	121.9	119.8	118.3
∠H2-Mn1-H1 (°)	82.2	80.5	82.1
τ_2 , Mn1	0.86	0.85	0.85
Mn1-Mn2 (Å)	2.871	2.871	2.814
∠Mn1-H1-Mn2 (°)	97.5	99.6	97.9
Spin-Densities ^a Mn1	4.91	4.75	-
Spin-Densities ^a Mn2	4.91	-4.75	-
Mayer Bond Order for Mn-Mn	0.21	0.27	-
Relative E_{ac}^b (kcal mol ⁻¹)	0.90	0	-

^aMulliken spin-density.
^bThe relative energies (kcal mol⁻¹) are referenced to the AF-coupled state.

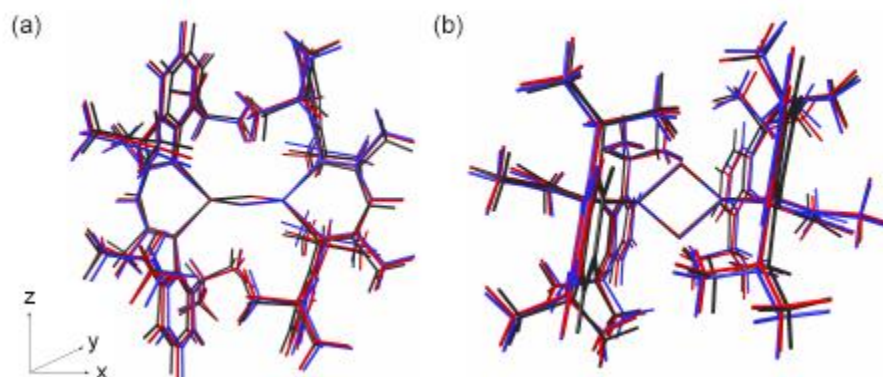


Fig. 2.3.3. Overlay of the structures of **1-CS** in black, **1-F** in blue, and **1-AF** in red: (a) the xz-plane view and (b) the orthogonal view along the z-axis.

To investigate the lack of a Mn–Mn multiple bond, the MO-diagrams of **1-F** and **1-AF**, shown in Fig. 2.3.4, were analyzed. The MO-diagrams were simplified by adopting a local D_{2h} symmetry. In total, 12 MOs comprising 10 metal-based orbitals and 2 hydride-based orbitals are presented. As expected, the in-phase combination of the d_{xy} orbitals interacts with the out-of-phase combination of

hydride orbitals to afford a bonding and antibonding combination for both **1-F** and **1-AF**. Furthermore, an interaction between the bonding combination of the d_{z^2} orbitals and the in-phase combination of hydride orbitals is observed, generating a low-lying bonding orbital and high-lying antibonding orbital.

The left-hand side of Fig. 2.3.5. illustrates the **1-F** MOs. The β -orbitals of **1-F** are higher in energy than the corresponding α -orbitals due to the absence of exchange interaction.^{80,81} On the other hand, both spin orbitals are paired up in **1-AF**, with MOs derived from Noodleman's broken symmetry approach, shown on the right-hand side of Fig. 2.3.4. Interestingly, the metal d_{xz} orbitals are higher in energy than expected. Isosurface plots of the in-phase and out-of-phase combinations of the d_{xz} orbitals are shown in Fig. 2.3.5a for **1-F** and 2.4.5b for **1-AF**. The metal d_{xz} orbitals are mixed with ligand-based orbitals via σ -bonding. The ligand-based orbitals are antibonding with respect to the metal-based orbitals, destabilizing the MOs. The corresponding bonding combinations are much lower in energy, and not depicted.

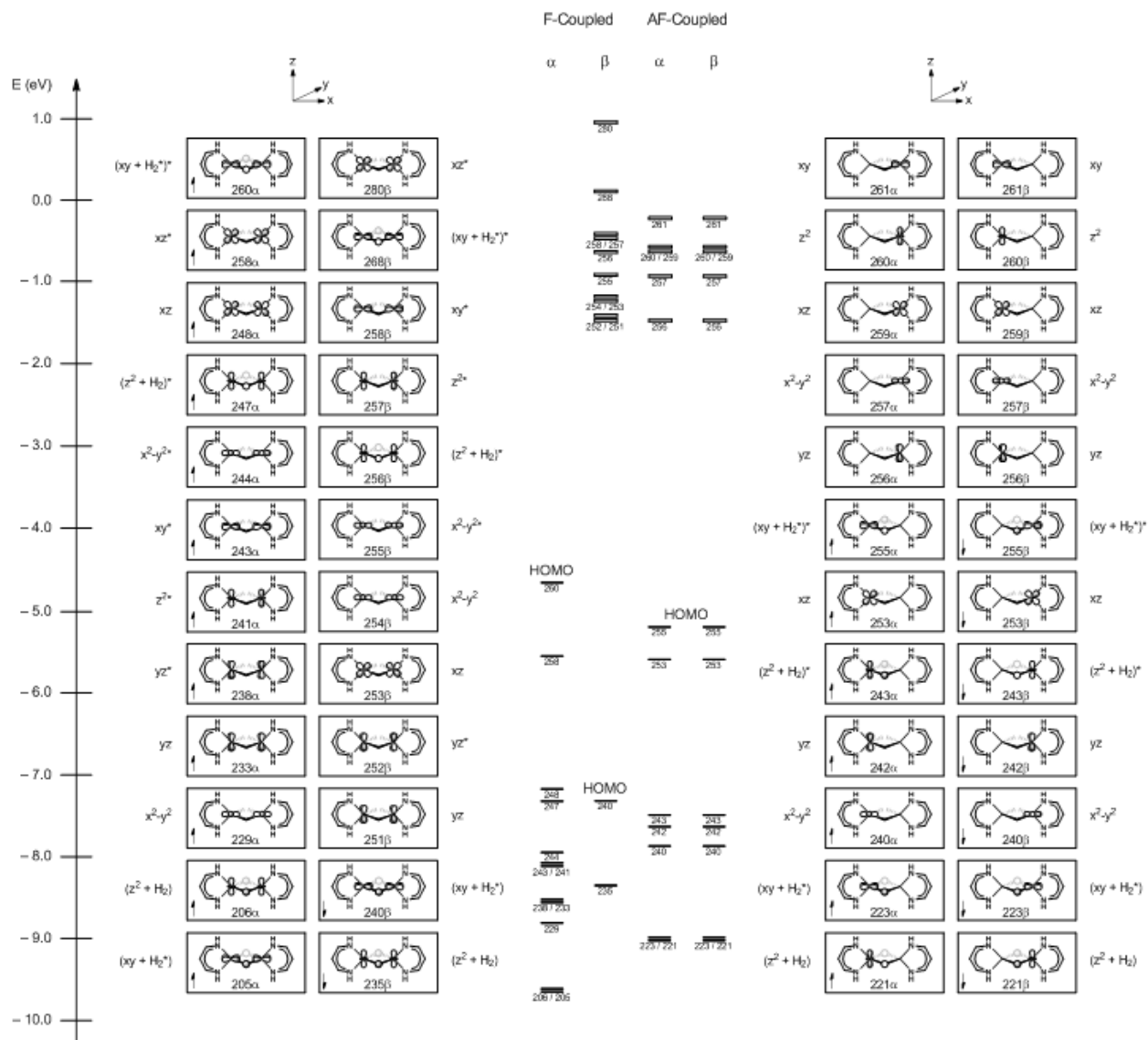


Fig. 2.3.4. MO-diagrams of ferro- (1-F) and antiferromagnetically-coupled (1-AF) Mn centers in 1.

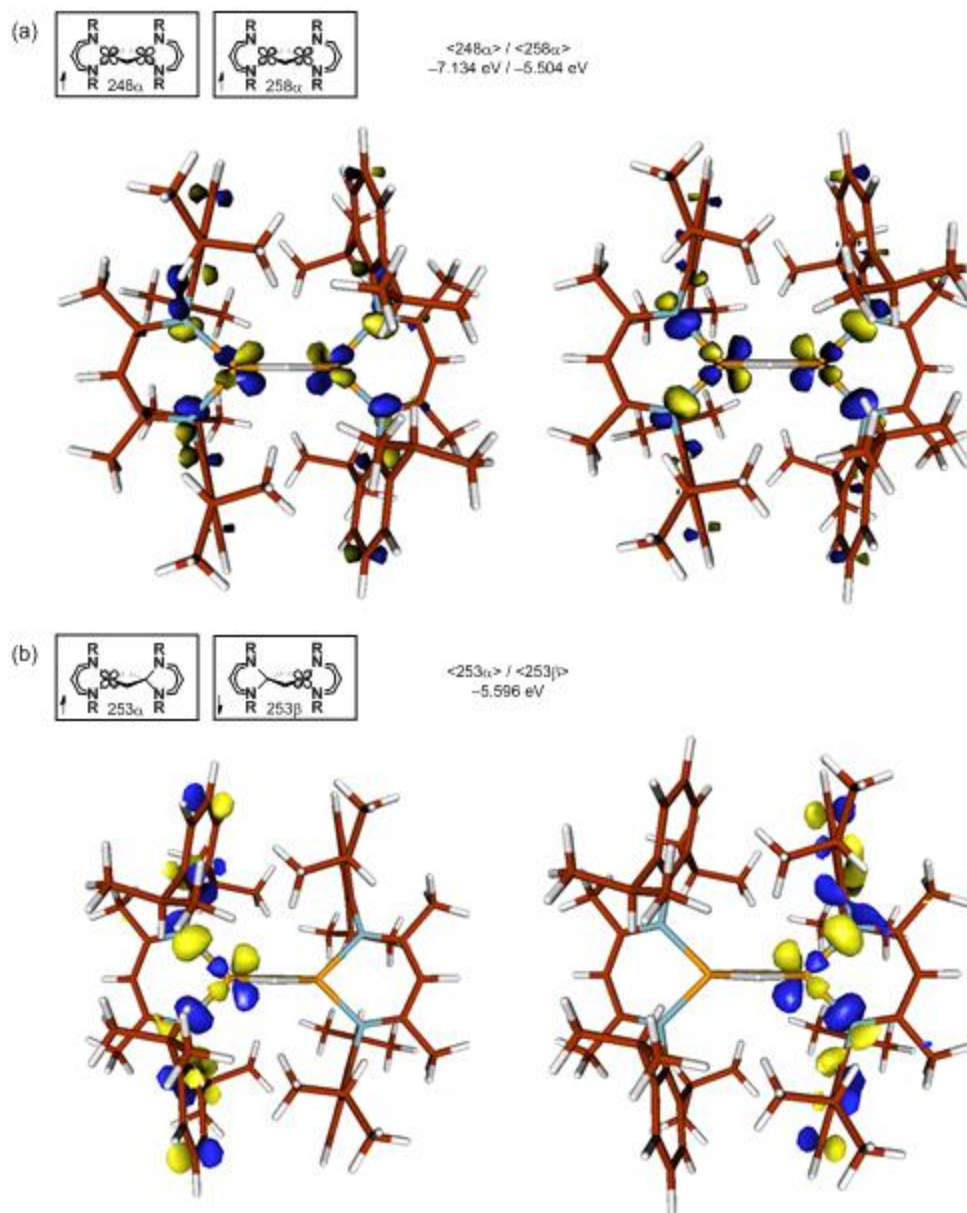


Fig. 2.3.5. Isosurface plots (isodensity = 0.05 au) of the metal d_{xz} -based MOs of **1**: a) metal–metal in-phase, and out-of-phase α -spin, **1-F** b) α -spin and β -spin, **1-AF**.

2.3.2 Determination of J_0 for **1**

According to Noodleman's description, the Heisenberg exchange coupling constant, J_0 , can be directly calculated from the energy difference between a high-spin state and a broken symmetry state.⁷⁴ In the broken-symmetry approach, the energy difference between the high spin state and the broken symmetry state is:

$$E(S_{\max}) - E(BS) = -S_{\max}(S_{\max} + 1) J_0 \quad (2.6)$$

where S_{\max} is the maximum number of unpaired electrons per monomeric unit. Here, each manganese center has five unpaired d -electrons; therefore, the high spin state is a sextet. The coupling constant is sensitive to the distance between metal centers, as well as the number of electrons shared by the metal centers. Therefore, probing J_o values will be the best way to assess the interaction between metal centers. Using equation (2.6) the Heisenberg exchange coupling constant was determined to be $J_o = -10.9 \text{ cm}^{-1}$. This is lower by a factor 10 than the J_o -value determined by Roesky and co-workers, at -110 cm^{-1} using the B3LYP functional,⁶² in which a Mn–Mn single bond is formed by the Mn-based $4s$ -orbitals. Jones and co-workers synthesized a similar compound, $(L^*Mn)_2$ ($L^* = -N(AR^*)SiMe_3$), which also shows a single Mn–Mn bond arising from σ -interactions.⁸² The calculated J_o -value was reported by the Jones' group to be -47.5 cm^{-1} , which is still significantly stronger than the one calculated here.

The Broken Symmetry approach does not distinguish between direct magnetic exchange interactions and indirect magnetic exchange interactions, such as superexchange. Bridging hydrides are capable of mediating strong superexchange interactions between metal centers due to the small size of the ion and the symmetry of the s -orbital. Calculating superexchange explicitly with DFT requires the use of a multi-reference calculation, such as Complete Active Space SCF (CASSCF),⁸³ which is computationally expensive for a complex of the size of **1**. Nevertheless, it is expected that the superexchange is at least partially responsible for the calculated J_o value.

With a computed J_o value in hand, attempts were made to determine this value experimentally using two complementary techniques by Professor Trovitch of Arizona State University and Professor John Anderson of the University of Chicago. In the initial contribution by Trovitch and coworkers, **1** was found to exhibit an ambient temperature magnetic moment of $5.2 \mu_B$ at 298 K (Evans Method); however, variable temperature data had not been collected.¹ Therefore, a sample of **1** was analyzed by SQUID magnetometry between 1.8–300 K by Anderson and coworkers. At room temperature the χT value is $4.5 \text{ cm}^3\text{K/mol}$ ($6\mu_B$), which is significantly below the spin-only value for two noninteracting $S = 5/2$ Mn(II) centers ($\chi T_{so} = 8.75$), indicating antiferromagnetic coupling between the Mn centers, a multiply-bonded complex with a lower overall spin state, or some intermediate case. Upon cooling, the χT of **1** decreases slightly, indicative of weak antiferromagnetic coupling (Fig. 2.3.6).

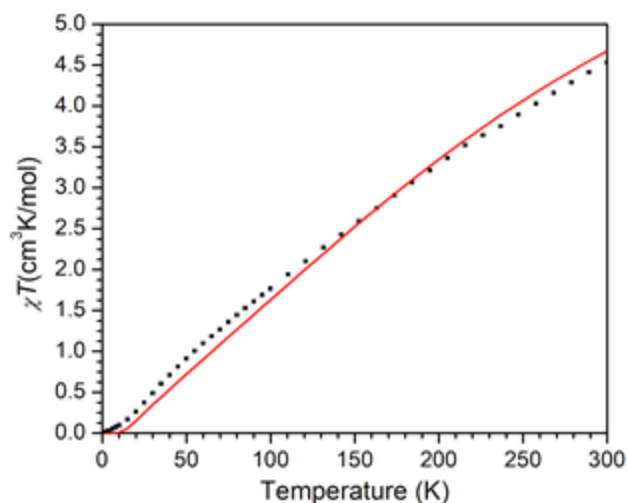


Fig. 2.3.6. Temperature-dependent χT vs. T data for **1** collected under an applied field of 0.1 T. The red line is the fit where g is fixed at 2.05, as outlined in the text. The measurements were performed by Professor Anderson and coworkers.

The data collected at 0.1 T was modeled using either the DAVE or PHI programs.^{84,85} The two Mn centers of **1** were treated as isotropic, spin-only ($S_{Mn} = 5/2$) centers and presumed to be identical given their crystallographic equivalence. The best fit was acquired with a g -value of 1.82 (± 0.25) and an antiferromagnetic J_o coupling of -15 cm^{-1} ($\pm 0.1 \text{ cm}^{-1}$, normalized to $2J_o$ convention). Constraining the fit to the g value of 2.05 previously determined by EPR spectroscopy resulted in a slightly poorer but still reasonable fit with $J_o = -20 \text{ cm}^{-1}$ ($\pm 0.1 \text{ cm}^{-1}$). The value of g is within error in both fits, and either fit provides a consistent J_o value of -15 to -20 cm^{-1} .

To determine the isotropic exchange coupling, J_o , by X-band EPR spectroscopy, Trovitch and coworkers have carried out measurements of **1** at six different temperatures between 4 K and 106 K. Fig. 2.3.7a shows the temperature variation of the EPR intensity of the signals occurring around 180 mT (B_1) and 60 mT (B_2) (see Fig. 2.3.7b). The EPR signals at B_1 and B_2 belong to the $S = 1$ and $S = 2$ spin manifolds, respectively. They were identified by simultaneously simulating the EPR spectra of **1** at different temperatures. At 4 K (Fig. 2.3.7b), the spectrum was simulated considering a superposition of two spin manifolds, $S = 1$ and $S = 2$. The parameters used to fit the EPR spectra of **1** (see EPR spectroscopy section) were within experimental error of those reported earlier by Mukhopadhyay et al.⁸ The value of J_o can be obtained by fitting the experimental data in Fig. 2.3.7a using equation (2.13) (see EPR spectroscopy section in the appendix), which accounts for the Boltzmann populations of the energy levels associated with the $S = 0, S = 1, S = 2, S = 3, S = 4$ and $S = 5$ spin manifolds. A value of -9.5 cm^{-1} was obtained when all data points were included in the fit. If

the data point corresponding to the EPR signal at B_1 and measured at 80 K is not included in the fit (see figure caption for details), a value of -10.8 cm^{-1} was obtained. Using the average of the two values, Trovitch and coworkers determined from the temperature dependence of the X-band EPR spectrum of **1** that $J_o = -10.2 \pm 0.7 \text{ cm}^{-1}$. Taken both SQUID and EPR data, our computational and the experimental efforts of Professors Trovitch and Anderson to determine J_o are consistent with weak antiferromagnetic coupling between the Mn centers of **1**.

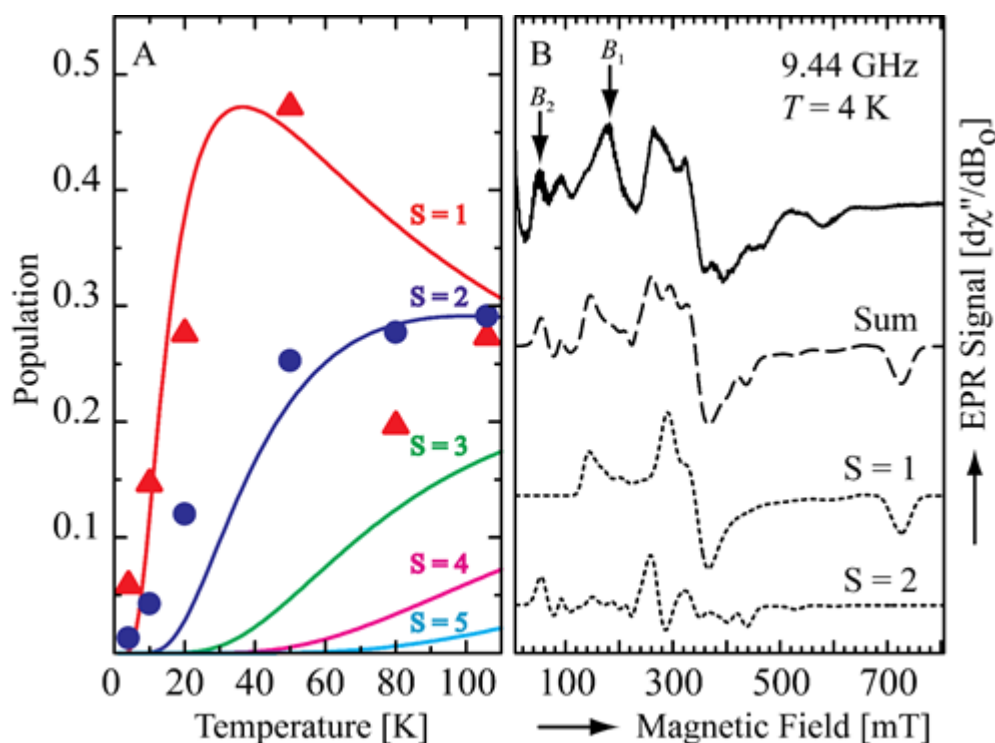


Fig. 2.3.7. a) Signal intensity times temperature as a function of temperature for the EPR signals around 180 mT (B_1) (red triangles) and 60 mT (B_2) (blue circles) present in the EPR spectra of **1**. Temperature dependence of Boltzman populations corresponding to $S = 1$ (red line), $S = 2$ (blue line), $S = 3$ (green line), $S = 4$ (magenta line) and $S = 5$ (cyan line) spin manifolds. These curves were calculated using Eq. 13 with the value of J_o that best fit the available experimental data. It can be noticed that the population corresponding to the EPR signal at B_1 and measured at 80 K deviates significantly from the fit. b) The X-band EPR spectrum of **1** at 4 K. The solid line is the experimental spectrum, and the dashed line is the sum of the simulated spectra for $S = 1$ and $S = 2$ dimer spin states (dotted lines). The EPR signals marked with arrows as B_1 and B_2 were identified as originating from the $S = 1$ and $S = 2$ dimer spin states, respectively. This spectrum did not show well-resolved ^{55}Mn hyperfine couplings. However, multiline patterns corresponding to the hyperfine couplings ($\sim 3.6 \text{ mT}$) of two equivalent Mn(II) ions were observed in the EPR spectrum at 20 K (not shown). The measurements were performed by Professor Trovitch and coworkers.

2.3.3 Influence of Steric Bulk

We initially hypothesized that the sterically bulky β -diketimate ligands preclude the formation of a metal–metal multiple bond as described in the literature.^{1,62,77,82} Bulky groups attached

to the β -diketiminate ligand may sterically prohibit the proximal positioning of two metal centers, despite the metal centers being within the range for interaction.⁸⁶ In this regard, the lower than anticipated bond order between the metal centers was proposed to be attributed to steric hindrance in bulky β -diketiminate ligands. To validate this proposition, a model compound supported by the most sterically simplified version of the β -diketiminate ligand, of which all substituents are replaced with hydrogen atoms, was modeled. Fig. 2.3.8. depicts the model compound, **2**.

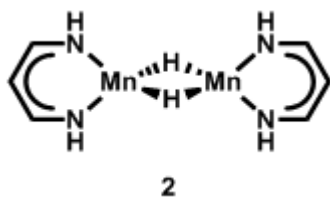
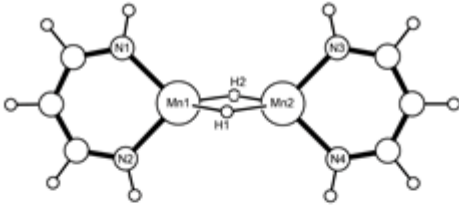


Fig. 2.3.8. The model compound **2**.

Calculations on **2** reproduced a nearly identical MO diagram to that of the full compound, **1** (Fig. S2.1). A structural comparison between the full and model compound is presented in Table 2.3.2. The Mn-Mn bond length of the model compound was calculated to be 2.816 and 2.789 Å for **2-F** and **2-AF**. The metal centers were moved apart from each other by only 0.06 Å and 0.08 Å for ferro- and antiferromagnetically coupled states, respectively. The corresponding Mayer bond orders are 0.21 and 0.24, which are comparable to the bond orders of 0.21 and 0.27 for **1-F** and **1-AF**. The steric effect, therefore, is small enough for **1-F** and **1-AF** to retain their bond orders.

Table 2.3.2 Comparison between optimized geometries of the full complex, **1**, and the model compound **2**.


	1-F	1-AF	2-F	2-AF
Spin-Coupling	F	AF	F	AF
Mn1-H1 (Å)	1.897	1.834	1.889	1.847
Mn1-H2 (Å)	1.919	1.923	1.894	1.856
Mn1-N1 (Å)	2.099	2.103	2.069	2.067
Mn1-N2 (Å)	2.116	2.110	2.069	2.067
∠H1-Mn1-N2 (°)	116.9	120.5	121.9	121.6
∠N2-Mn1-N1 (°)	90.1	90.5	91.4	91.5
∠N1-Mn1-H2 (°)	121.9	119.8	120.7	121.1
∠H2-Mn1-H1 (°)	82.2	80.5	83.7	82.3
α , Mn1	0.86	0.85	0.83	0.83
Mn1-Mn2 (Å)	2.871	2.871	2.816	2.789
∠Mn1-H1-Mn2 (°)	97.5	99.6	96.2	97.7
Spin-Densities* Mn1	4.91	4.75	4.87	4.73
Spin-Densities* Mn2	4.91	-4.75	4.87	-4.73
Mayer Bond Order for Mn-Mn	0.21	0.27	0.21	0.24
Relative E_{2z^2} (kcal mol ⁻¹)	0.90	0	4.24	0

*Mulliken spin-density.
^aThe relative energies (kcal mol⁻¹) are referenced to the AF-coupled states.

The irrelevance of steric hindrance caused by the ligand to the bond order between the metal centers led us to consider the fundamental basis for the absence of a multiple bond between the Mn centers. Interaction between the metal d-orbitals and ligand-based orbitals may be an intrinsic reason for the unexpectedly low Mn-Mn bond order. The β -diketiminato ligand has a bite angle of around 90° (Table 2.3.1 and 2.3.2) and a π -conjugated system, which affords thorough mixing between the metal- and ligand-based orbitals.

As previously shown for **1**, the d_{xz} orbitals are mixed with the orbitals of the nitrogen atoms of the β -diketiminato ligands (Fig. S2.5). The metal d_{xy} orbitals also interact with the β -diketiminato-based orbitals (Fig. 2.3.9a and c), as well as with the hydride-based orbitals (Fig. 2.3.9b and d). Electrons supposed to occupy bonding orbitals between the metal centers are dispersed to metal-ligand interactions, resulting in the attenuation of the metal-metal bond.^{86,87}

Despite adopting the broken symmetry approach, the d_{xy} and d_{z^2} orbitals contain delocalized counterparts (Fig. 2.3.9c and d, Fig. 2.3.10c and d). Considering weak δ -interactions with the d_{z^2} orbitals, we posited that the Mn-Mn bond order around 0.2 is mainly attributed to the d_{xy} orbitals. The MO diagram for the ferromagnetically coupled state, **2-F**, corroborates this postulation (Fig. S2.1). The α -orbitals consist of pairs of orbitals composed of a bonding combination and the corresponding antibonding combination, which compensate each other. The remaining bonding combinations in the β -orbitals coincide with the d_{xy} and d_{z^2} orbitals in the antiferromagnetically

coupled state, which contain residual moieties on the Mn-center on which an MO is not localized, as shown in Fig. 2.3.9c and d and Fig. 2.3.10c and d. Furthermore, hydride orbitals are incorporated in the d_{xy} and the d_{z^2} orbitals, presumably indicating the presence of a superexchange interaction mediated by bridging hydrides.

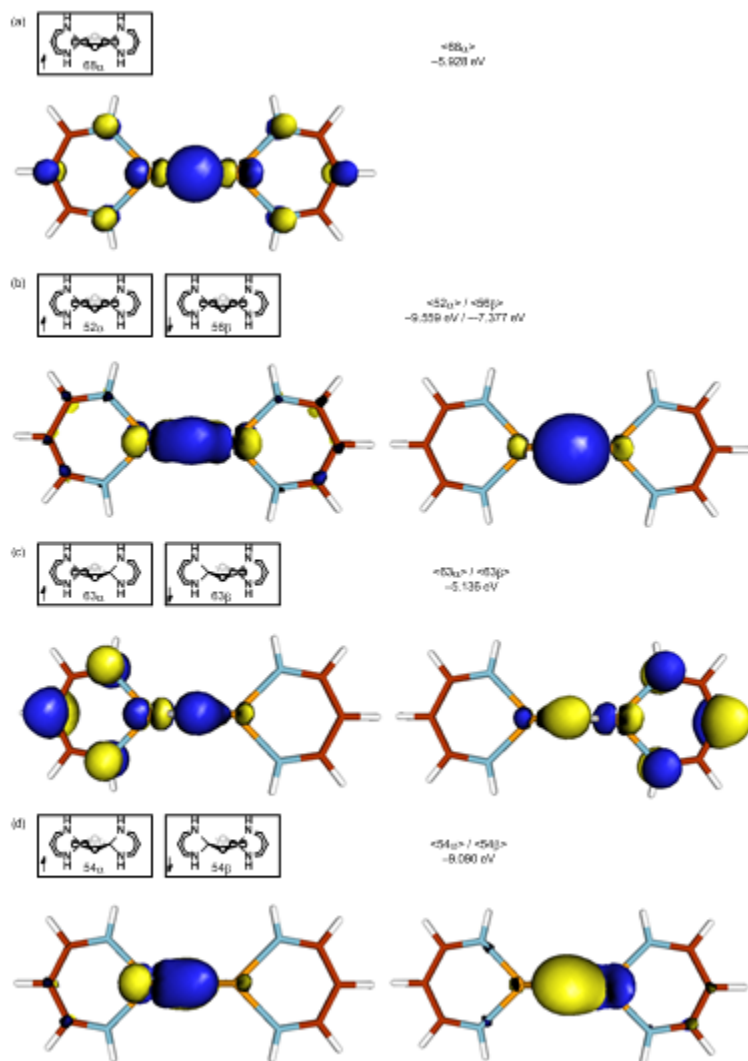


Fig. 2.3.9. Isosurface plots (isodensity = 0.05 au) of the metal d_{xy} -based MOs of **2**: a) **2-F**, $\langle 68\alpha \rangle$, b) **2-F**, $\langle 52\alpha \rangle$ and $\langle 56\beta \rangle$, c) **2-AF**, $\langle 63\alpha \rangle$ and $\langle 63\beta \rangle$, and d) **2-AF**, $\langle 54\alpha \rangle$ and $\langle 54\beta \rangle$.

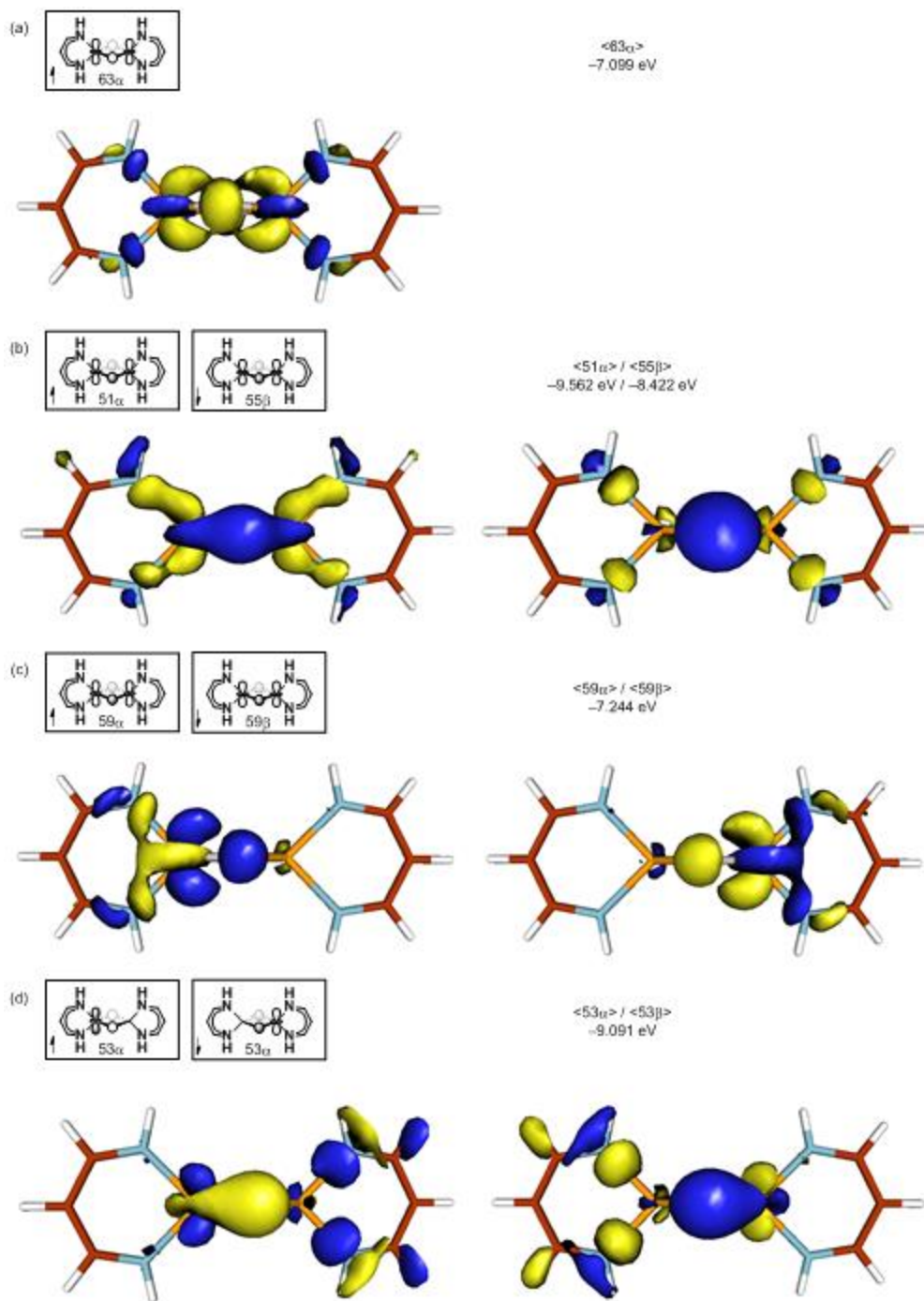


Fig. 2.3.10. Isosurface plots (isodensity = 0.05 au) of the metal d_z -based MOs of **2**: a) **2-F**, $\langle 63\alpha \rangle$, b) **2-F**, $\langle 51\alpha \rangle$ and $\langle 55\beta \rangle$, c) **2-AF**, $\langle 59\alpha \rangle$ and $\langle 59\beta \rangle$, and d) **2-AF**, $\langle 53\alpha \rangle$ and $\langle 53\beta \rangle$.

2.3.4 Extension to Ligands with *p*-orbitals

Catalytic nitrile dihydroboration with **1** was recently reported by the Trovitch group.⁹⁴ The proposed mechanism involves an Mn(II) intermediate similar to **1**, in which bridging hydrides are replaced with bridging imino groups. This intermediate is far more stable than **1** by 22.51 kcal/mol. The Mn–Mn distance is also elongated from 2.814 Å to 3.082 Å. The dissimilarity was proposed to be attributed to a fundamental difference in orbitals. The bonding combination between the metal centers in **2** interacts through bridging ligands. Spherical 1s-orbitals in hydrides enable the in-phase combination of metal *d*-orbitals, while *p*-orbitals having a node will preclude mediation of the bonding interaction between the metal centers. The out-of-phase combination will solely be allowed with *p*-orbitals.

To evaluate this proposition, interactions between the Mn centers were further investigated by substituting bridging ligands in which *p*-orbitals are available, fluorides and hydroxides, for the bridging hydrides, of which the respective model compounds **3** and **4** are shown in Fig. 2.3.11. The structural disparity induced by the availability of *p*-orbitals is highlighted in Table 2.4.3. The most remarkable difference is the elongated Mn–Mn distance from 2.789 Å to 3.074 Å and 3.100 Å in conjunction with the decrease in the Mayer bond order from 0.24 to 0.02 and 0.03 in **3** and **4**, respectively. Similar patterns for mixing metal- and BDI ligand-based orbitals were observed, except for the incorporation of *p*-orbitals. All metal *d*-orbitals were mixed with *p*-orbitals, implying that the degree of the Mn–Mn interaction is decreased. The *d_{xy}* orbitals assigned as the main factor to the bonding interaction between the metal centers are depicted in Fig. 2.3.12 and Fig. 2.3.13. The *p_x* and *p_y* orbitals are mixed to interact with the *d_{xy}* orbitals. Bonding interactions between the metal centers via bridging ligands are symmetrically forbidden as we anticipated. As a consequence, the Mn–Mn bond strength is attenuated, resulting in the increase in Mn–Mn distance.

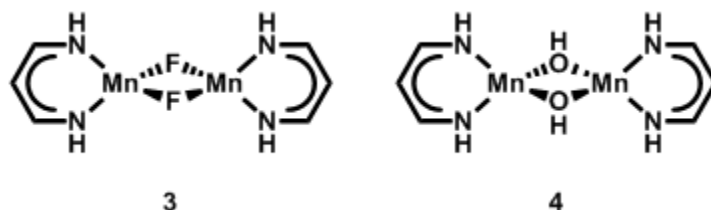


Fig. 2.3.11. The model compound with bridging fluorides, **3**, and the model compound with bridging hydroxides, **4**.

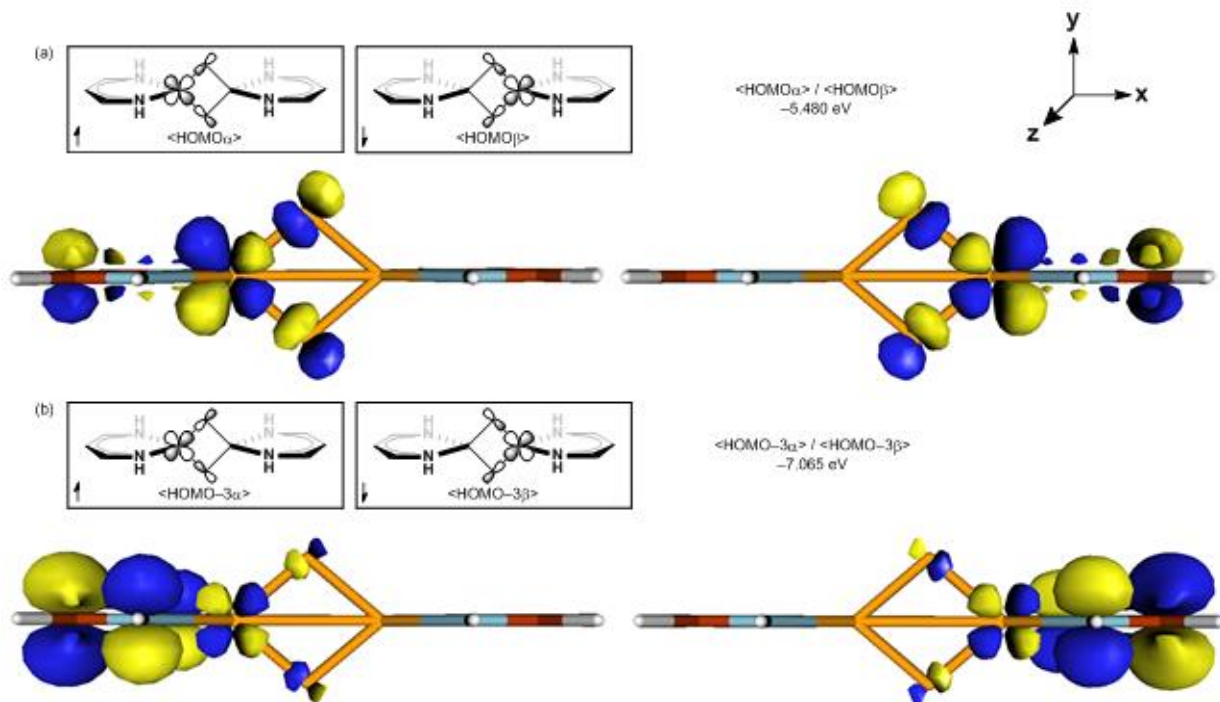


Fig. 2.3.12. Isosurface plots (isodensity = 0.05 au) of the metal d_{xy} -based MOs of **3**: (a) HOMO, (b) HOMO-3. To clearly visualize directions of orbitals, the Cartesian coordinate system consistently adopted is rotated 90 degrees in counterclockwise along the X-axis, as shown in the upper right corner.

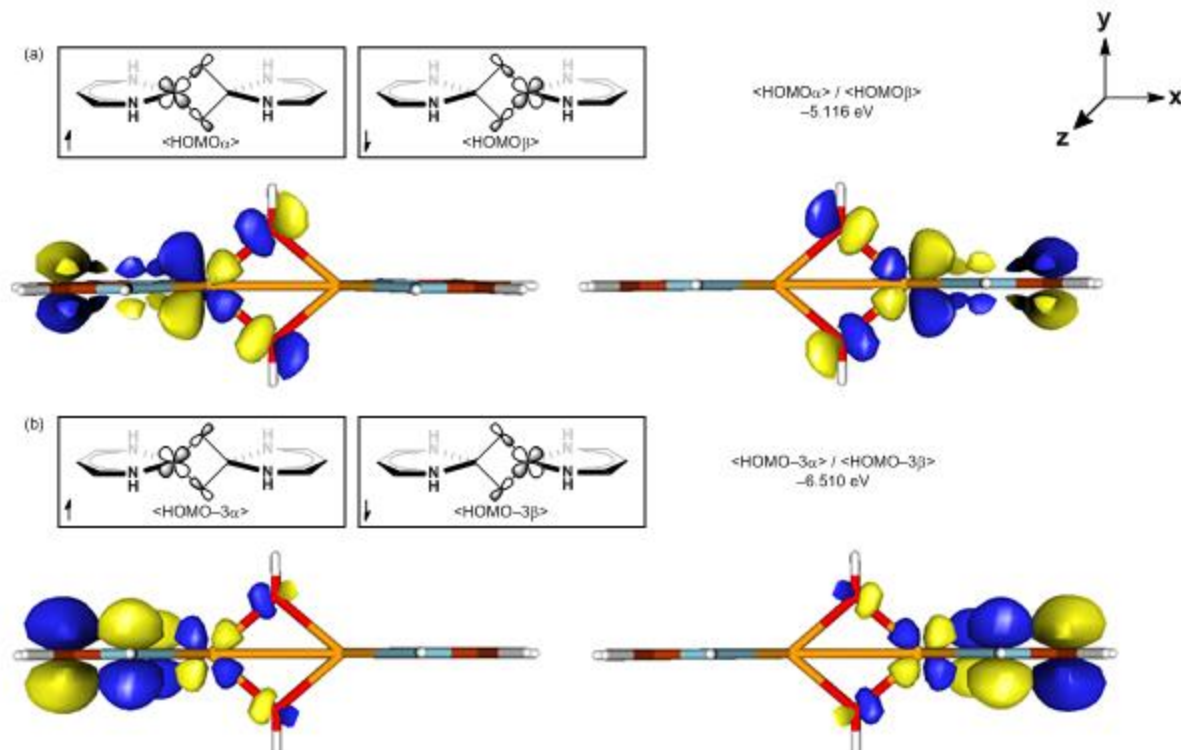


Fig. 2.3.13. Isosurface plots (isodensity = 0.05 au) of the metal d_{xy} -based MOs of **4**: (a) HOMO, (b) HOMO-3. To clearly visualize directions of orbitals, the Cartesian coordinate system consistently adopted is rotated 90 degrees in counterclockwise by the X-axis and shown at the upper right corner.

Table 2.3.3 Comparison of structural parameters of the optimized geometries of the model compound **2**, **3**, and **4**

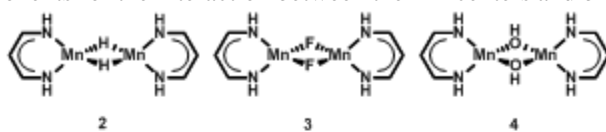
	2	3	4
Spin-Coupling	AF	AF	AF
Mn1-H1(O1, F1) (Å)	1.847	2.019	2.043
Mn1-H2(O2, F2) (Å)	1.856	2.019	2.043
Mn1-N1 (Å)	2.067	2.076	2.101
Mn1-N2 (Å)	2.067	2.077	2.102
\angle H1(O1, F1)-Mn1-N2 ($^\circ$)	121.6	121.8	130.0
\angle N2-Mn1-N1 ($^\circ$)	91.5	92.4	90.6
\angle N1-Mn1-H2(O2, F2) ($^\circ$)	121.1	121.7	122.6
\angle H2(O2, F2)-Mn1-H1(O1, F1) ($^\circ$)	82.3	80.8	81.3
τ , Mn1	0.83	0.82	0.76
Mn1-Mn2 (Å)	2.789	3.074	3.100
\angle Mn1-H1(O1, F1)-Mn2 ($^\circ$)	97.7	99.2	98.7
Spin-Densities ^a Mn1	4.73	4.81	4.80
Spin-Densities ^a Mn2	-4.73	-4.81	-4.80
Mayer Bond Order for Mn-Mn	0.24	0.02	0.03

^aMulliken spin-density.

The lengthening of the distance between the metal centers presumably accounts for the stabilization effect with p -orbitals. Table 2.3.4 describes the fragment analysis of the three model

compounds, **2**, **3**, and **4**. The compounds were fragmented into three parts, one manganese subunit, the other manganese subunit and the bridging ligands. Large electron densities located on the bridging anions give rise to an immense repulsion in the core composed of the Mn centers and bridging ligands. Due to the enlarged core structure of **3**, the Pauli repulsion term is decreased by 22.5 kcal/mol compared to **2**. On the contrary, the elongated Mn–Mn bond decreases the extent of electrostatic interaction in **3**; however, the electrostatic terms do not exceed the repulsion terms, rendering **3** more stable than **2**. At first glance, the computed interaction components of **4** seem contradictory, but they can be explained. The size of hydroxides is much larger than that of hydrides or fluorides; therefore, the Pauli repulsion term of **4** is 10.0 kcal/mol higher than that of **2** due to proximity. Electrostatic interactions are understandably strengthened. Consequently, degrees of stabilization through non-orbital interactions in **3** and **4** are nearly identical. Despite the longer distances, magnitudes of orbital interactions in **3** and **4** are greater than that in **2**. The participation of *p*-orbitals probably contributes to the greater orbital interactions. The penalties in non-orbital interaction and orbital interactions of **2** gives rise to electronic destabilization. Dimerization involves a stabilization in enthalpy and disadvantage in entropy. The penalties in interaction terms of **2** can partially offset the stabilizing effect in enthalpy of the dimerization. As a result, the electronic energy of the dimer is lower by 32.5 kcal/mol than that of two equivalents of the corresponding monomer in the case of **2**, while **3** and **4** show greater energy differences of 41.9 and 50.0 kcal/mol, respectively. Taking into account entropy penalties, the monomer-dimer Gibbs energy differences in **2**, **3**, and **4** at room temperature diminish to 18.6, 28.3 and 35.8 kcal/mol, respectively. Furthermore, in benzene solution, the energy difference at room temperature of **2** becomes 4.1 kcal/mol, while that of **3** and **4** stays at 16.5 and 24.6 kcal/mol, which is far more difficult to attain. In accordance with the model study, **1-AF**, the fluoride congener of **1**, $[(2,6\text{-iPr}_2\text{PhBDI})\text{Mn}(\mu\text{-F})_2]$ (**5**), and the hydroxide congener of **1**, $[(2,6\text{-iPr}_2\text{PhBDI})\text{Mn}(\mu\text{-OH})_2]$ (**6**) are lower in energy than their corresponding monomers by 7.1, 22.6 and 28.8 kcal/mol, respectively.

Table 2.3.4. Computed components for the interaction between the Mn centers and bridging ligands in **2**, **3** and **4**.



(In kcal/mol)	2	3	4	$\Delta_1(3-2)$	$\Delta_2(4-2)$
Mn–Mn (Å)	2.789	3.074	3.100	–	–
Pauli repulsion (PR)	278.93	256.47	288.91	–22.46	9.98
Electrostatic Interaction (EI)	–618.85	–608.69	–639.48	10.16	–20.63
Non-orbital Interaction (PR + EI)	–339.92	–352.22	–350.58	–12.30	–10.65
Orbital Interaction	–395.21	–400.01	–431.05	–4.80	–35.84

Synthesis of the Hydroxide-Bridged Dimer.¹ To support the calculations conducted on model compound **4** and full hydroxide complex **6**, attempts to synthesize the latter were made by Trovitch and coworkers. Upon dissolving **1** in THF solution, the slow addition of a second THF solution containing two equivalents of H₂O resulted in the liberation of H₂ gas and a light yellow product identified as **6**. The magnetic susceptibility of **6** was analyzed by the Gouy method and determined to be 7.4 μ_B at 291 K. Notably, the infrared spectrum of this complex was found to exhibit a hydroxide OH stretch at 3,695 cm^{–1} that shifted to 2,726 cm^{–1} when D₂O was employed in the synthesis (**6-d₂**).

To confirm the identity of **6** and obtain an experimental Mn–Mn bond distance, single crystals were grown from THF at 238 K and X-ray diffraction data was collected. The crystal structure of **6** (Fig. 2.3.14) features an inversion center, eclipsed BDI ligands, and a relatively long Mn1–Mn1A distance of 3.1426(9) Å. This distance is consistent with the calculated distance of 3.189 Å and a full comparison of the experimental and computational bond distances is provided in Table 2.3.5.

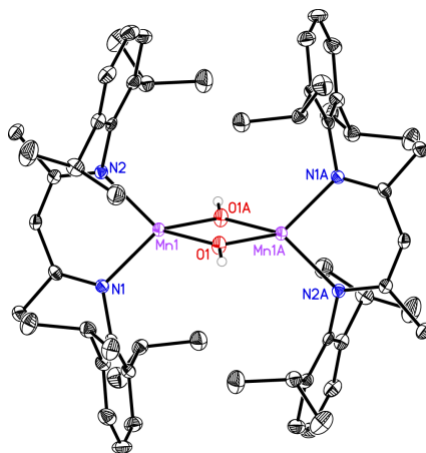
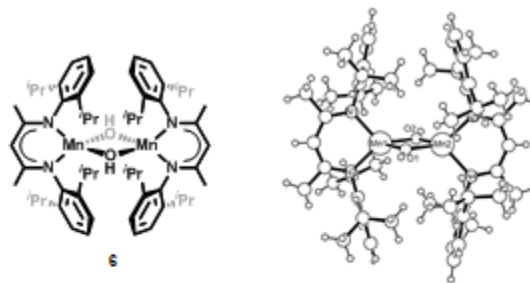


Fig. 2.3.14. Solid state structure of **6** at 30% probability ellipsoids. From Professor Trovitch and coworkers.

¹ The synthesis and characterization of this compound were carried out by Professor Trovitch and coworkers.

Table 2.3.5. Experimental and calculated metrical parameters of **6**



	6 exp.	6 calc.
Spin-Coupling	-	AF
Mn1-O1 (Å)	2.037	2.056
Mn1-O2 (Å)	2.054	2.071
Mn1-N1 (Å)	2.106	2.107
Mn1-N2 (Å)	2.108	2.119
∠O1-Mn1-N2 (°)	127.0	125.0
∠N2-Mn1-N1 (°)	91.3	89.2
∠N1-Mn1-O2 (°)	120.5	124.5
∠O2-Mn1-O1 (°)	79.6	78.9
τ_4 Mn1	0.80	0.78
Mn1-Mn2 (Å)	3.142	3.189
∠Mn1-O1-Mn2 (°)	100.4	101.2
Spin-Densities* Mn1	-	4.80
Spin-Densities* Mn2	-	-4.80

*Mulliken spin-density.

2.4 Conclusion

DFT calculations were performed on the ferromagnetically and antiferromagnetically coupled manganese dimer, $[(2,6\text{-iPr}_2\text{PhBDI})\text{Mn}(\mu\text{-H})_2]$, in order to elucidate interactions between the metal centers. Whereas the electron count suggests a formal bond order of 3, the calculated bond order is much lower at 0.21 and 0.27, for **1-F** and **1-AF** respectively. The total energies of **1-F** and **1-AF** show that the antiferromagnetically coupled complex is the ground state, for which a coupling value $J_o = -10.9 \text{ cm}^{-1}$ was calculated. This is in accordance with experimental data, which afforded $J_o = -15 \pm 0.1 \text{ cm}^{-1}$ or $J_o = -20 \pm 0.1 \text{ cm}^{-1}$ by SQUID and $J_o = -10.2 \pm 0.7 \text{ cm}^{-1}$ by EPR spectroscopy. The influence of steric bulk on the decreased bond order was ruled out by calculations on model compounds **2-F** and **2-AF**, which show near identical bond orders to **1-F** and **1-AF**. Analysis of the molecular orbitals reveals that extensive mixing of the metal d_{xy} - and d_{xz} -orbitals with the ligand-based orbitals prevents direct orbital overlap of the metal centers, reducing the bond order. The d_{xy} orbitals are likely the major contributor to bonding interaction due to the presence of residual moiety in the broken-symmetry orbitals. A stable analogue of **1** comprising an increased Mn-Mn distance reported by the Trovitch group led us to further investigate the effect of bridging p -orbitals. The p -orbitals

symmetrically prohibit the in-phase combination of metal *d*-orbitals as well as interact with all metal *d*-orbitals, resulting in the lengthening of the Mn–Mn bond. Fragment analysis revealed that both non-orbital interactions and orbital interactions stabilize dimeric compounds with bridging ligands with *p*-orbitals. To corroborate our calculations, the congener of **1** with bridging hydroxides was synthesized and characterized.

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Appendix

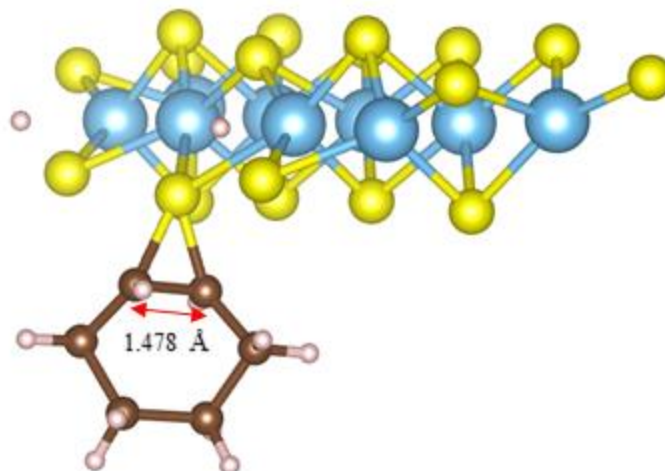


Figure S1.1. Thiirane intermediate of cyclohexane. The C–C bond of the carbon atoms participating in the thiirane cycle is shortened slightly from 1.498 Å to 1.478 Å, but is longer than the typical C = C bond at 1.340 Å. The total energy of the complex $\Delta E_{\text{Tot}} = 41.16$ kcal/mol.

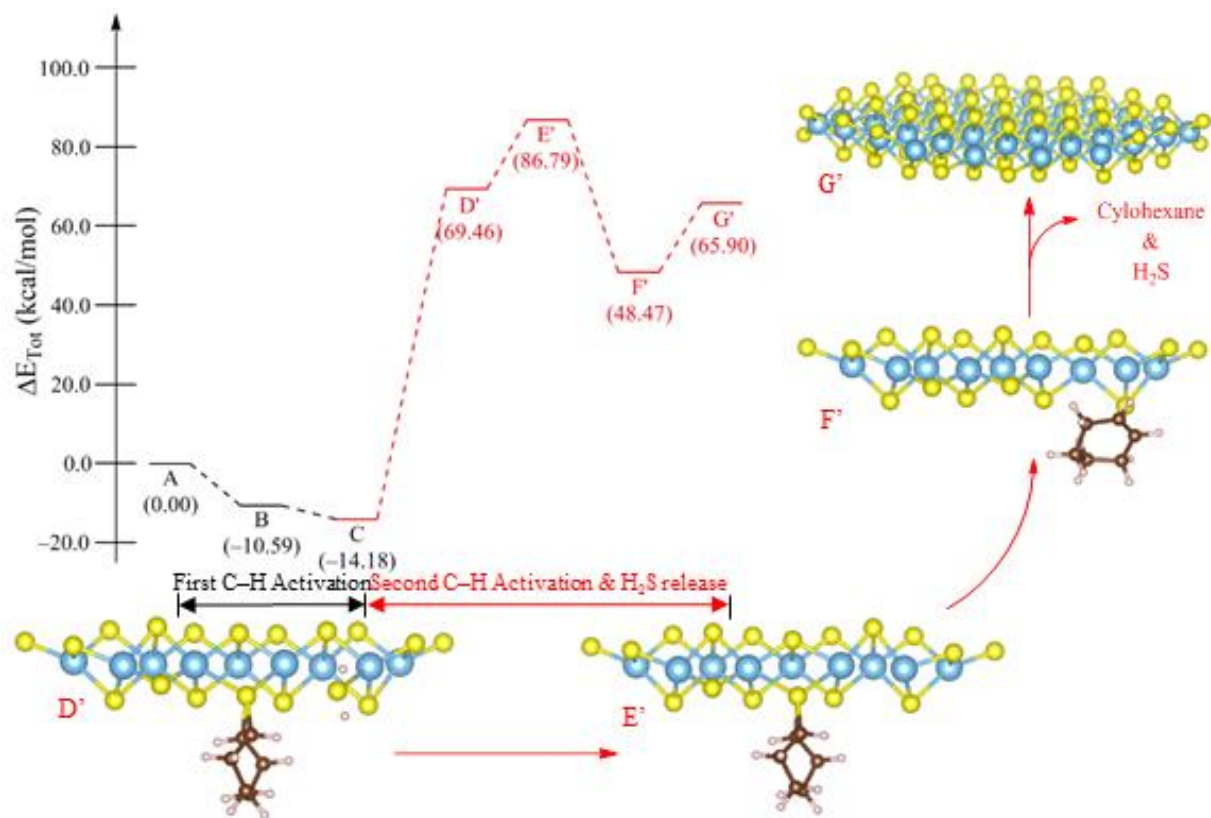


Figure S1.2. The reaction energy profile of cyclohexane dehydrogenation including H₂S production. The part of the profile given in red is the part that differs from the reaction energy profile given in Fig. 1.3.1. Only the geometries of the differing reaction steps are shown. ΔE_{Tot} given in kcal/mol.

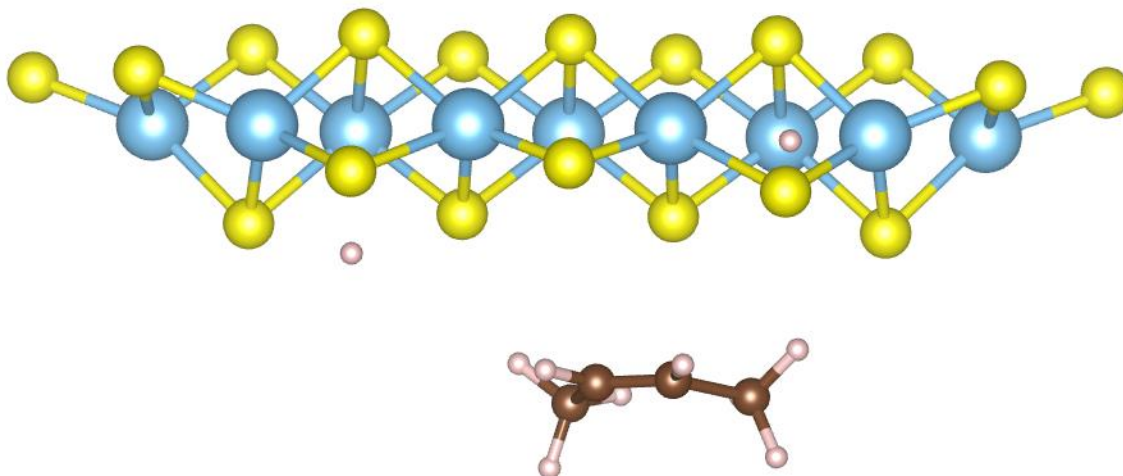


Figure S1.3. Geometry of 5', which was excluded from the energy diagram in Fig. 3.x due to the direct release of 2-butene without formation of a thirane intermediate. $\Delta E_{\text{Tot}} = 8.89$ kcal/mol, which is 0.77 kcal/mol higher than that of 4'.

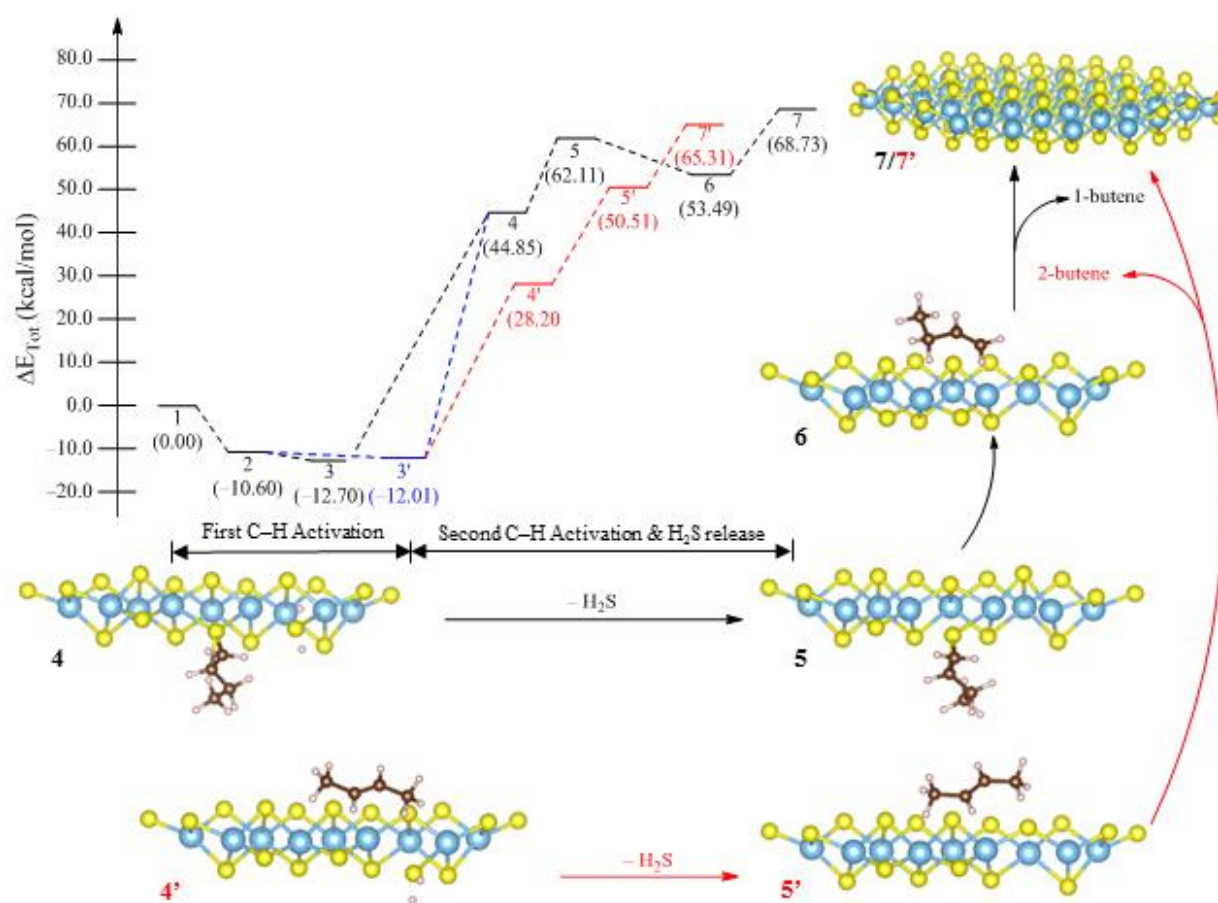


Figure S1.4. The reaction energy profile of butane dehydrogenation including H_2S production. Steps 4–7 and 4'–7' differ from the reaction energy profile given in Fig. 1.4.3. Only the geometries of the differing reaction steps are shown. ΔE_{Tot} given in kcal/mol.

Table S1.1. Energies of the optimized geometries

Compound name	Energy in kcal/mol
Cyclohexane Dehydrogenation	
A	-199969.4814
Cyclohexane	-12039.48167
B	-212019.5494
C	-212023.1428
D	-212006.5362
E	-200650.5883
Cyclohexene	-11338.75824
H2S removed mechanism	
D'	-211939.501
E'	-208929.2972
F'	-208967.6193
G'	-197611.4315
H2S	-2992.873591
Cyclohexane thiirane intermediate	-211967.808

Butane Dehydrogenation

1	-199969.4814
butane	-8699.179545
2	-208679.2637
3	-208681.3596
4	-208651.5919
5	-208657.5181
6/6'	-200650.5883
1-butene	-7995.621202
3'	-208680.6692
4'	-208660.5366
2-butene	-7999.050036
H2S removed Mechanism	
4	-208623.8075
5	-205613.6794
6	-205622.2992
7/7'	-197611.4315
4'	-208640.4613

5'	-205625.2816
H2S	-2992.873591
dissociated thiirane intermediate	-208659.7686
2-methylthietane	-208318.1831
tetrahydrothiophene	-208329.1519

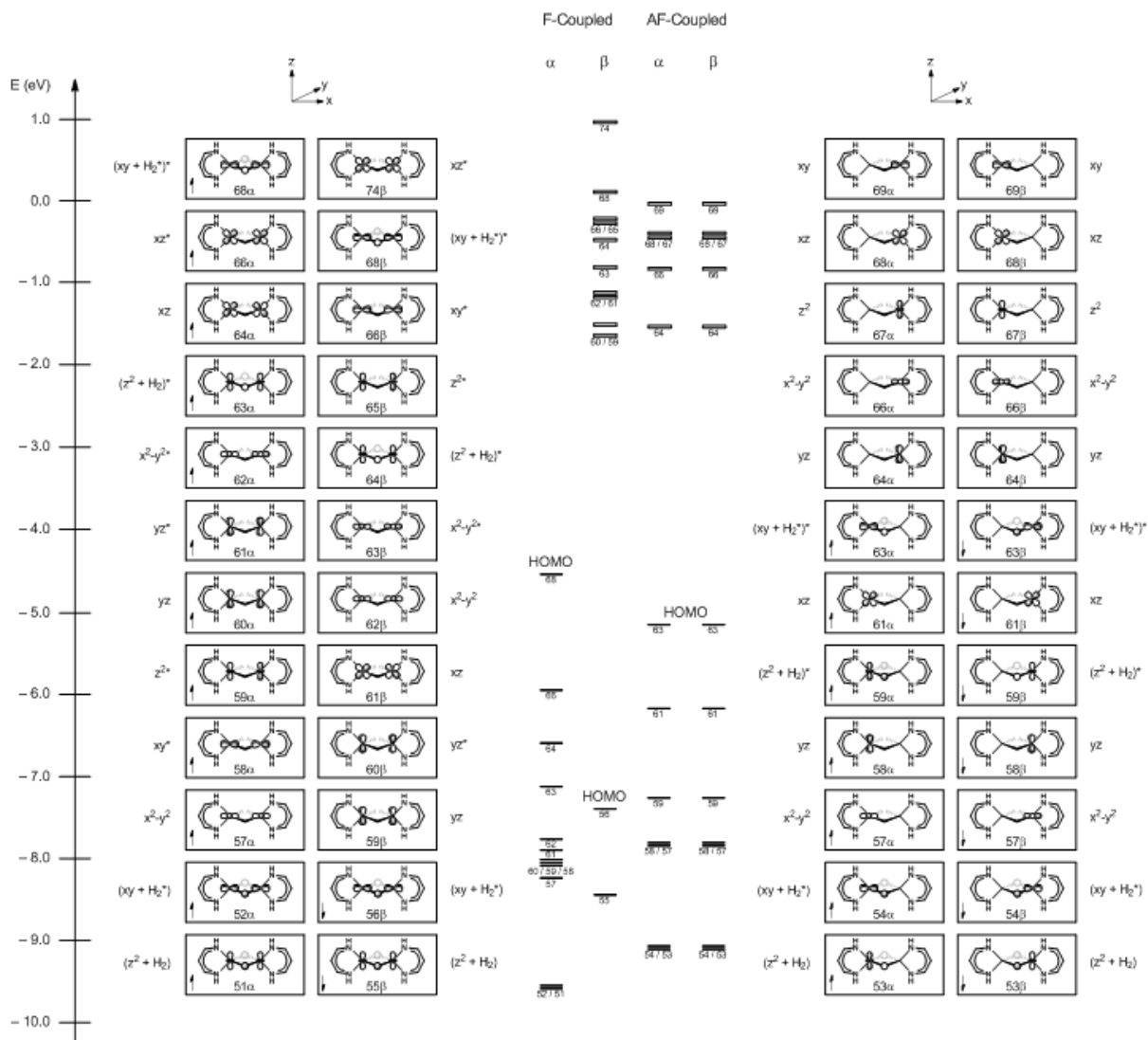


Figure S2.1. The MO-diagrams of ferro (2-F)- and antiferromagnetically (2-AF) coupled manganese centers of the model compound.

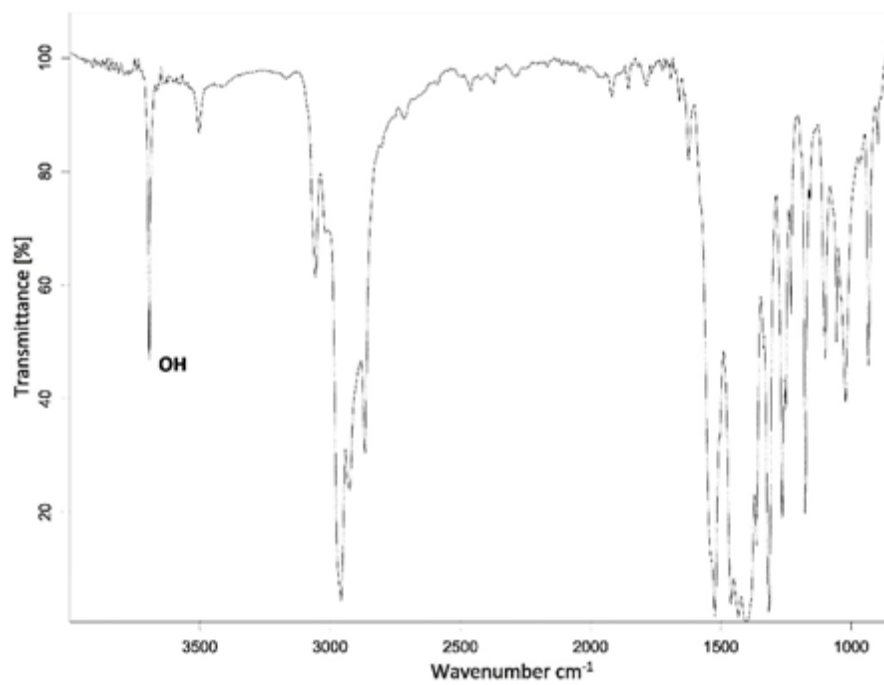


Figure S2.2. Solid-state infrared spectrum of **4** in KBr.

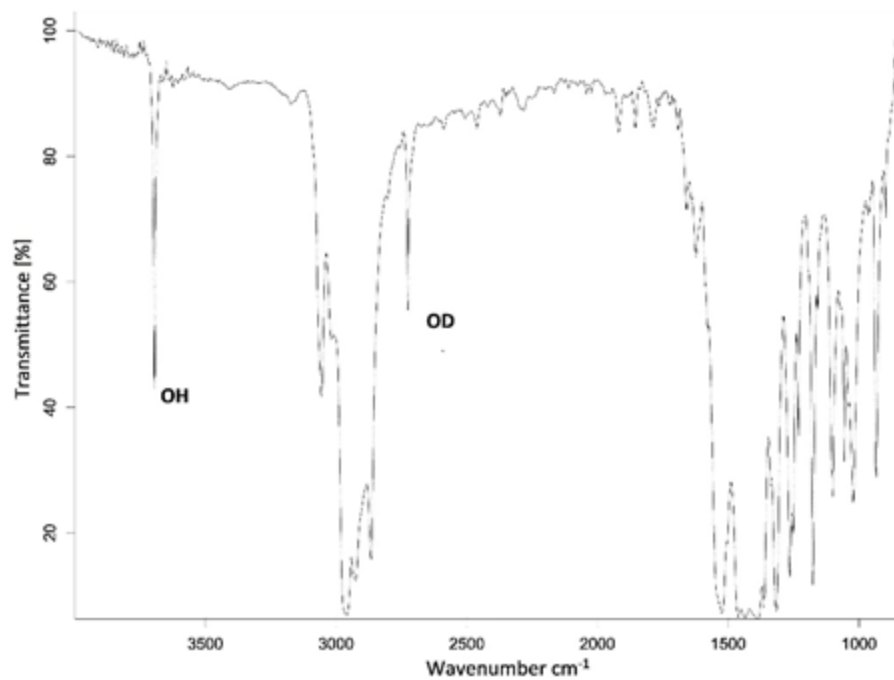


Figure S2.3. Solid-state infrared spectrum of **4-*d*₂** in KBr.

CRYSTALLOGRAPHIC DATA

Table S2.1. Crystallographic data for **4**.

	4
chemical formula	C ₅₈ H ₈₄ Mn ₂ N ₄ O ₂
formula weight	979.17
crystal dimensions	0.193 x 0.134 x 0.073
crystal system	monoclinic
space group	C 1 2/c 1
<i>a</i> (Å)	22.969(3)
<i>b</i> (Å)	14.6534(17)
<i>c</i> (Å)	16.1143(19)
α (deg)	90
β (deg)	91.558(2)
γ (deg)	90
<i>V</i> (Å ³)	5421.6(11)
<i>Z</i>	4
<i>T</i> (°C)	123.00(10)
ρ_{calcd} (g cm ⁻³)	1.200
μ (mm ⁻¹)	0.509

reflections collected	22387
data/restraints/parameters	4978/0/312
R ₁ [I > 2σ(I)]	0.0560
wR ₂ (all data)	0.01464
Goodness-of-fit	1.028
Largest peak, hole (eÅ ⁻³)	0.800, -0.294

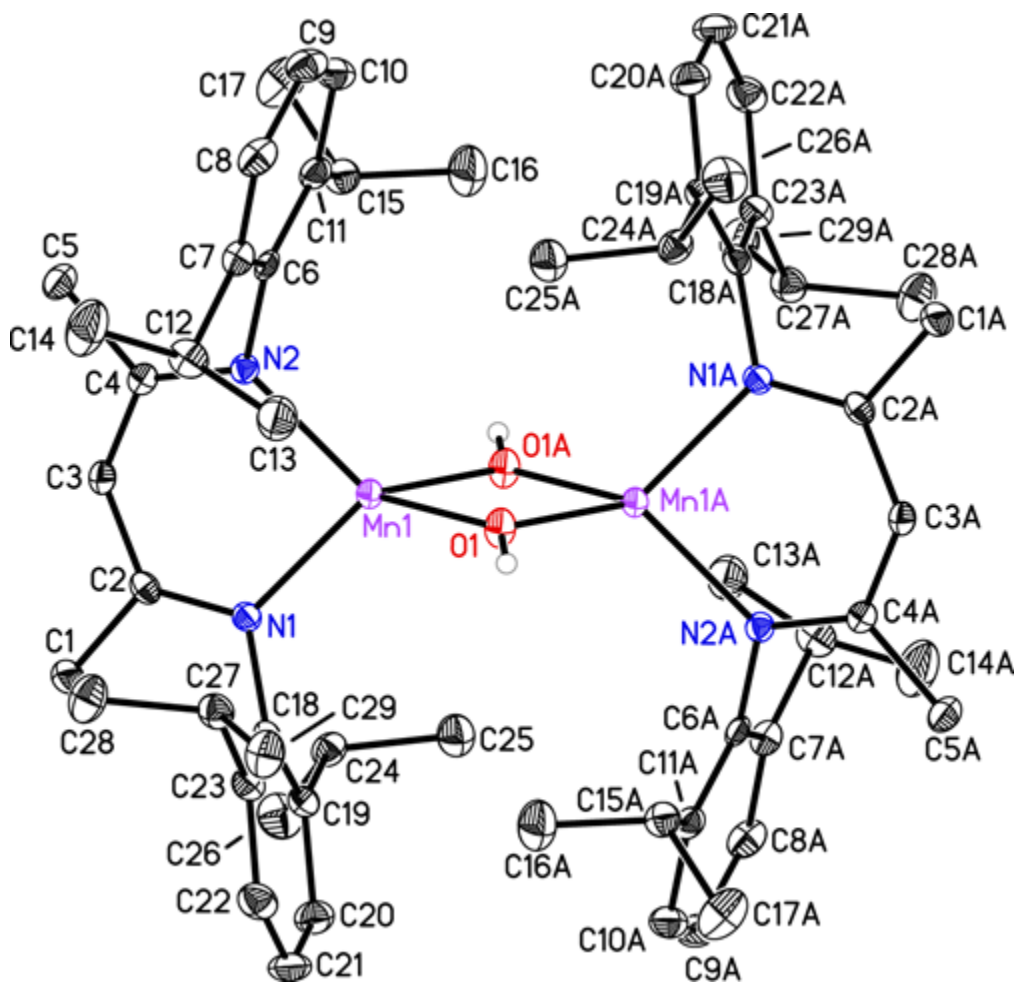


Figure S2.4. The molecular structure of **6** shown at 30% probability ellipsoids. Hydrogen atoms other than the hydroxyl hydrogens are omitted for clarity.

Table S2.2. Bond distances (Å) and angles (°) for **6**.

Mn1-O1	2.038(2)	C3-C4	1.404(4)	C15-C17	1.529(4)
Mn1-O1A	2.054(2)	C4-C5	1.514(4)	C18-C19	1.401(4)
Mn1-N2	2.106(2)	C6-C11	1.401(4)	C18-C23	1.403(4)
Mn1-N1	2.108(2)	C6-C7	1.411(4)	C19-C20	1.391(4)
Mn1-Mn1A	3.1426(9)	C7-C8	1.382(4)	C19-C24	1.511(4)
O1-Mn1A	2.054(2)	C7-C12	1.516(4)	C20-C21	1.382(5)
O1-H1	0.71(4)	C8-C9	1.382(5)	C21-C22	1.381(5)
N1-C2	1.338(4)	C9-C10	1.374(5)	C22-C23	1.396(4)
N1-C18	1.442(4)	C10-C11	1.404(4)	C23-C27	1.526(4)
N2-C4	1.333(4)	C11-C15	1.509(4)	C24-C25	1.523(4)
N2-C6	1.444(3)	C12-C13	1.497(5)	C24-C26	1.527(4)
C1-C2	1.509(4)	C12-C14	1.526(5)	C27-C29	1.520(4)
C2-C3	1.398(4)	C15-C16	1.510(5)	C27-C28	1.528(4)
O1-Mn1-O1A	79.64(10)	C2-C3-C4	130.5(3)	C16-C15-C17	109.8(3)
O1-Mn1-N2	126.96(9)	N2-C4-C3	124.5(3)	C19-C18-C23	121.8(3)
O1A-Mn1-N2	120.82(9)	N2-C4-C5	120.2(2)	C19-C18-N1	120.5(3)
O1-Mn1-N1	121.34(9)	C3-C4-C5	115.2(2)	C23-C18-N1	117.7(3)
N2-Mn1-N1	91.34(9)	C11-C6-C7	121.0(3)	C20-C19-C18	117.6(3)
O1-Mn1-Mn1A	40.01(6)	C11-C6-N2	120.0(3)	C20-C19-C24	120.4(3)
O1A-Mn1-Mn1A	39.63(6)	C7-C6-N2	119.0(2)	C18-C19-C24	122.0(3)
N2-Mn1-Mn1A	136.44(7)	C8-C7-C6	118.3(3)	C21-C20-C19	121.6(3)
N1-Mn1-Mn1A	132.01(7)	C8-C7-C12	120.7(3)	C22-C21-C20	120.1(3)
Mn1-O1-Mn1A	100.37(10)	C6-C7-C12	121.0(3)	C21-C22-C23	120.6(3)
Mn1-O1-H1	134.(3)	C9-C8-C7	121.7(3)	C22-C23-C18	118.2(3)
Mn1A-O1-H1	120.(3)	C10-C9-C8	119.7(3)	C22-C23-C27	120.4(3)
C2-N1-C18	116.5(2)	C9-C10-C11	121.3(3)	C18-C23-C27	121.4(3)
C2-N1-Mn1	121.27(18)	C6-C11-C10	118.1(3)	C19-C24-C25	111.1(3)
C18-N1-Mn1	122.24(17)	C6-C11-C15	122.0(3)	C19-C24-C26	113.4(3)
C4-N2-C6	116.5(2)	C10-C11-C15	119.8(3)	C25-C24-C26	109.3(3)
C4-N2-Mn1	121.75(19)	C13-C12-C7	110.7(3)	C29-C27-C23	113.3(3)
C6-N2-Mn1	121.69(17)	C13-C12-C14	111.2(3)	C29-C27-C28	110.2(3)
N1-C2-C3	124.5(3)	C7-C12-C14	113.3(3)	C23-C27-C28	111.4(3)
N1-C2-C1	120.1(3)	C11-C15-C16	112.1(3)		
C3-C2-C1	115.4(3)	C11-C15-C17	111.6(3)		

Table S2.3 Cartesian coordinates for the optimized geometries.

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l-F

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2-F

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3-AF

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3-MONO

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4-AF

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4-MONO

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5-AF

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5-MONO

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=====

6-AF

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H 0.826724648 2.686411619 18.320426941
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H -3.599677086 0.964053988 12.883426666
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H -3.456509352 -1.744072318 9.854543686
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H -0.101620466 0.700492680 11.774504662
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6-MONO

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H -1.908984065 -2.610122204 18.215118408
H -0.772026002 -1.252243996 18.379161835
C -1.083516002 -1.889251232 16.340402603
C -1.150066018 -3.192414045 15.811742783
H -1.538879037 -3.950246096 16.480497360
C -0.802978039 -3.639234781 14.520083427
C -1.023731589 -5.105853081 14.204523087
H -0.080256715 -5.583872795 13.920087814
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H -1.695782781 -5.211623192 13.346236229
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C -1.729110837 1.332426906 15.963718414
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C 1.708125710 0.022655040 17.128265381
H 1.365266442 -1.009925961 17.014183044
C 2.337988853 0.146607876 18.524517059
H 1.590872765 -0.000300122 19.311428070

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C 0.665350139 -4.214112282 9.696271896
H 0.924385488 -4.541289330 8.692951202
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H 2.584352493 -4.722275734 10.511997223
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C 2.335996628 -3.937807322 13.133926392
H 1.808396459 -3.905486822 14.091876030
C 3.225044727 -2.679469585 13.062407494

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H 3.837654114 -5.269437313 12.242552757
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H -2.403704643 -2.436995983 12.561883926
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H -3.470344305 -3.699387550 9.979516983
H -3.420285463 -4.490669727 11.561082840
H -4.433168411 -3.055640936 11.314092636
C -2.393916368 -1.227623582 10.794611931
H -3.349562645 -0.742385387 11.027252197
H -1.582504153 -0.565995693 11.118410110
H -2.327666759 -1.334724545 9.705533981
O 0.258020401 0.485319227 12.404191017
H 0.351406425 1.427744865 12.578169823

Table S2.4. Energies of the optimized geometries

	G(sol)
1-AF	-73136.623
1-F	-73136.655
2-AF	-18039.239
2-F	-18039.185
3-AF	-23446.852
3-MONO	-11723.068
4-AF	-22138.135
4-MONO	-11068.534
5-AF	-78544.608
5-MONO	-39271.814
6-AF	-76152.524

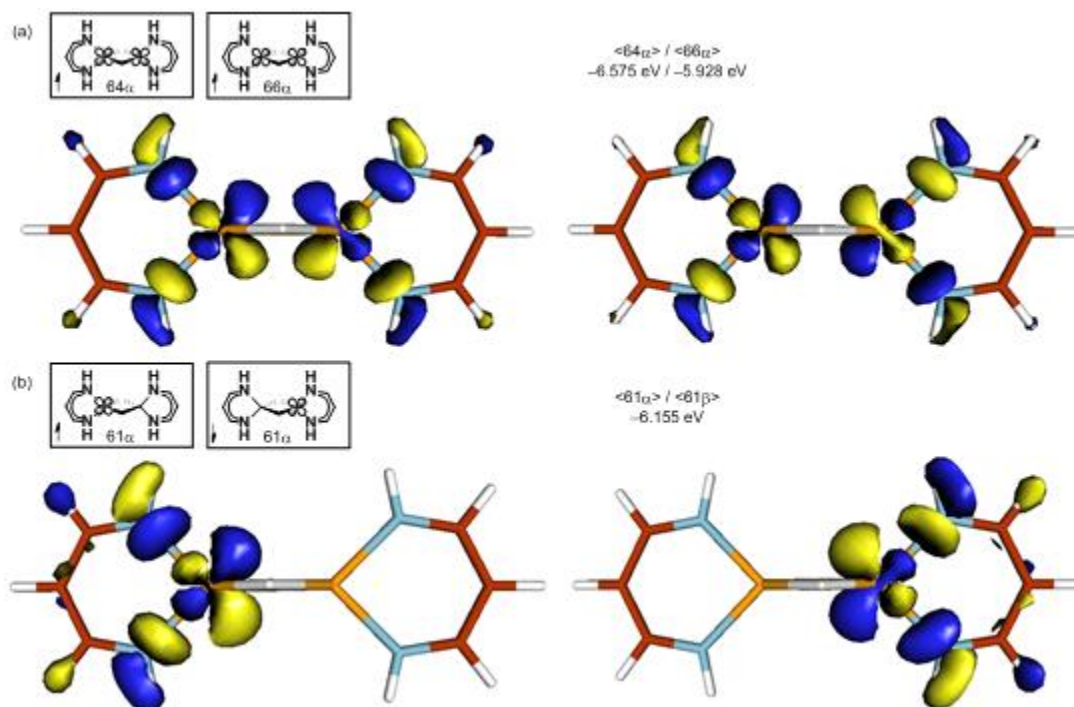


Fig. S2.5. Isosurface plots (isodensity = 0.05 au) of the metal d_{xz} -based MOs of **2**: metal–metal in-phase, and out-of-phase α -spin, **2-F** (b) α -spin and β -spin, **2-AF**.